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THE GAUSS-WEINGARTEN FORMULAE FOR THE HOMOGENEOUS LIFT TO THE OSCULATOR BUNDLE OF A FINSLER METRIC

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Abstract

In this article we present a study of the subspaces of the manifold OscM, the total space of the osculator bundle of a real manifold M. We obtain the induced connections of the canonical metrical N-linear connection determined by the homogeneous prolongation of a Finsler metric to the manifold OscM. We present the Gauss-Weingarten equations of the associated osculator submanifold.

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1 Introduction

The Sasaki N-prolongation \mathbb{G} to the osculator bundle without the null section $\widetilde{OscM} = OscM \setminus \{0\}$ of a Finslerian metric g_{ab} on the manifold M given by

$$\mathbb{G} = g_{ab}(x, y) \, dx^a \otimes dx^b + g_{ab}(x, y) \, \delta y^a \otimes \delta y^b \tag{*}$$

is a Riemannian structure on OscM, which depends only on the metric g_{ab} .

The tensor \mathbb{G} is not invariant with respect to the homothetis on the fibres of \widetilde{OscM} , because \mathbb{G} is not homogeneous with respect to the variable y^a .

In this paper, we use a new kind of prolongation \mathring{G} to OscM, ([8]), which depends only on the metric g_{ab} . Thus, \mathring{G} determines on the manifold OscM a Riemannian structure which is 0-homogeneous on the fibres of OscM.

Some geometrical properties of \mathbb{G} are studied: the canonical metrical N-linear connection, the induced linear connections etc.

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2 Preliminaries

As far as we know the general theory of submanifolds (in particular the Finsler submanifolds or the complex Finsler submanifolds) is far from being settled ([1],[10], [3],[11], [12]). In [9] and [10] R.Miron and M. Anastasiei give the theory of subspaces in generalized Lagrange spaces. Also, in [6] and [5] R. Miron presented the theory of subspaces in higher order Finsler and Lagrange spaces respectively.

If M is an immersed manifold in manifold M, a nonlinear connection on OscM induces a nonlinear connection N on OscM.

The d-tensor \mathbb{G} from (*) is not homogeneous with respect to the variable y^a . This is an incovenient from the point of view of mechanics. Moreover, the physical dimensions of the terms of \mathbb{G} are not the same. This disadvantage was corrected by R. Miron. He took a new kind of prolongation $\overset{\circ}{\mathbb{G}}$ to OscM of the fundamental tensor of a Finsler space, ([8]) (5), which depends only on the metric g_{ab} . Thus, $\overset{\circ}{\mathbb{G}}$ determines on the manifold OscM a Riemannian structure which is 0-homogeneous on the fibres of OscM and p is a positive constant required by applications in order that the physical dimensions of the terms of $\overset{\circ}{\mathbb{G}}$ be the same. He proved that there exist metrical N-linear connections with respect to the metric tensor $\overset{\circ}{\mathbb{G}}$.

We take this canonical N-linear metric connection D on the manifold OscM and obtain the induced tangent and normal connections and the relative covariant derivation in the algebra of d-tensor fields ([13], [16]).

In this paper we get the Gauss-Weingarten formulae of submanifold $Osc\dot{M}$.

Let us consider $F^n = (M, F)$ a Finsler space ([10]), and $F : TM = OscM \to \mathbb{R}$ the fundamental function. F is a C^{∞} function on the manifold OscM and it is continuous on the null section of the projection $\pi : OscM \to M$. The fundamental tensor on F^n is

$$g_{ab}\left(x,y
ight)=rac{1}{2}rac{\partial^{2}F^{2}}{\partial y^{a}\partial y^{b}}, \ \forall\left(x,y
ight)\in OscM.$$

The lagrangian $F^2(x, y)$ determines the canonical spray $S = y^a \frac{\partial}{\partial x^a} - 2G^a \frac{\partial}{\partial y^a}$ with the coefficients $G^a = \frac{1}{2}\gamma^a_{bc}(x, y) y^b y^c$, where $\gamma^a_{bc}(x, y)$ are the Christoffels symbols of the metric tensor $g_{ab}(x, y)$. The Cartan nonlinear connection N of the space F^n has the coefficients

$$N^a{}_b = \frac{\partial G^a}{\partial y^b}.\tag{1}$$

N determines a distribution on the manifold OscM, ([10],[9]), which is supplementary to the vertical distribution V. We have the next decomposition

The adapted basis of this decomposition is $\left\{\frac{\delta}{\delta x^a}, \frac{\partial}{\partial y^a}\right\}$, (a = 1, .., n) and its dual basis is $(dx^a, \delta y^a)$, where

$$\begin{cases} \frac{\delta}{\delta x^{a}} = \frac{\partial}{\partial x^{a}} - N^{b}{}_{a}\frac{\delta}{\delta y^{b}}, \\ \frac{\partial}{\partial y^{a}} = \frac{\partial}{\partial y^{a}} \end{cases}$$
(3)

and

$$\begin{cases} dx^a = dx^a, \\ \delta y^a = dy^a + N^a{}_b dx^b. \end{cases}$$
(4)

We use the next notations:

$$\delta_a = \frac{\delta}{\delta x^a}, \ \dot{\partial}_{1a} = \frac{\partial}{\partial y^a}$$

The fundamental tensor g_{ab} determines on the manifold \widetilde{OscM} the homogeneous N-lift $\overset{0}{\mathbb{G}}$,[8],

$$\overset{0}{\mathbb{G}} = g_{ab}\left(x,y\right) dx^{a} \otimes dx^{b} + h_{ab}\left(x,y\right) \delta y^{a} \otimes \delta y^{b},\tag{5}$$

where

$$h_{ab}(x,y) = \frac{p^2}{\|y\|^2} g_{ab}(x,y), \qquad (6)$$
$$\|y\|^2 = g_{ab}(x,y) y^a y^b.$$

This is homogeneous with respect to y, and p is a positive constant required by applications in order that the physical dimensions of the terms of \mathbb{G} be the same.

Let \check{M} be a real, m-dimensional manifold, immersed in M through the immersion $i: \check{M} \to M$. Localy, i can be given in the form

$$x^{a} = x^{a} \left(u^{1}, ..., u^{m} \right), \qquad rank \left\| \frac{\partial x^{a}}{\partial u^{\alpha}} \right\| = m.$$

The indices a, b, c,...run over the set $\{1, ..., n\}$ and $\alpha, \beta, \gamma, ...$ run on the set $\{1, ..., m\}$. We assume 1 < m < n. We take the immersed submanifold OscM of the manifold OscM, by the immersion $Osci : OscM \to OscM$. The parametric equations of the submanifold OscM are

$$\begin{cases} x^{a} = x^{a} \left(u^{1}, ..., u^{m} \right), rang \left\| \frac{\partial x^{a}}{\partial u^{\alpha}} \right\| = m \\ y^{a} = \frac{\partial x^{a}}{\partial u^{\alpha}} v^{\alpha}. \end{cases}$$

$$(7)$$

The restriction of the fundamental function F to the submanifold OscM is

$$\check{F}(u,v) = F(x(u), y(u,v))$$

and we call $\check{F}^m = (\check{M}, \check{F})$ the **induced Finsler subspaces** of F^n and \check{F} the **induced fundamental function**.

induced fundamental function. Let $B^a_{\alpha}(u) = \frac{\partial x^a}{\partial u^{\alpha}}$ and $g_{\alpha\beta}$ the induced fundamental tensor,

$$g_{\alpha\beta}(u,v) = g_{ab}\left(x\left(u\right), y\left(u,v\right)\right) B^{a}_{\alpha}B^{b}_{\beta}.$$
(8)

We obtain a system of d-vectors $\{B^a_{\alpha}, B^a_{\bar{\alpha}}\}$ which determines a moving frame $\mathcal{R} = \{(u, v); B^a_{\alpha}(u), B^a_{\bar{\alpha}}(u, v)\}$ in OscM along to the submanifold $Osc\dot{M}$.

Its dual frame will be denoted by $\mathcal{R}^* = \{B_a^{\alpha}(u,v), B_a^{\bar{\alpha}}(u,v)\}$. This is also defined on an open set $\check{\pi}^{-1}(\check{U}) \subset Osc\check{M}, \check{U}$ being a domain of a local chart on the submanifold \check{M} .

The conditions of duality are given by:

$$B^{a}_{\beta}B^{\alpha}_{a} = \delta^{\alpha}_{\beta}, \quad B^{a}_{\beta}B^{\bar{\alpha}}_{a} = 0, \quad B^{\alpha}_{a}B^{a}_{\bar{\beta}} = 0, \quad B^{\bar{\alpha}}_{a}B^{a}_{\bar{\beta}} = \delta^{\bar{\alpha}}_{\bar{\beta}}$$
$$B^{a}_{\alpha}B^{\alpha}_{b} + B^{a}_{\bar{\alpha}}B^{\bar{\alpha}}_{b} = \delta^{a}_{b}.$$

The restriction of the nonlinear connection N to $Osc\check{M}$ uniquely determines an induced nonlinear connection \check{N} on $\widetilde{Osc\check{M}}$

$$\check{N}^{\alpha}{}_{\beta} = B^{\alpha}_{a} \left(B^{a}_{0\beta} + N^{a}{}_{b} B^{b}_{\beta} \right).$$
⁽⁹⁾

The cobasis $(dx^i, \delta y^a)$ restricted to $Osc\check{M}$ is uniquely represented in the moving frame \mathcal{R} in the following form:

$$\begin{cases} dx^{a} = B^{a}_{\beta} du^{\beta} \\ \delta y^{a} = B^{a}_{\alpha} \delta v^{\alpha} + B^{a}_{\bar{\alpha}} K^{\bar{\alpha}}_{\beta} du^{\beta} \end{cases}$$
(10)

where

$$K^{\bar{\alpha}}_{\beta} = B^{\bar{\alpha}}_a \left(B^a_{0\beta} + M^a_b B^b_{\beta} \right), \ B^a_{0\beta} = B^a_{\alpha\beta} v^a.$$

A linear connection D on the manifold OscM is called **metrical N-linear connection** with respect to $\mathring{\mathbb{G}}$, if $D\mathring{\mathbb{G}} = 0$ and D preserves by parallelism the distributions N and V. The coefficients of the N-linear connections $D\Gamma(N)$ will be denoted with $\begin{pmatrix} H & a \\ L & bc \\ (00) & bc \end{pmatrix}, \begin{pmatrix} H & a \\ L & bc \\ (10) & bc \end{pmatrix}, \begin{pmatrix} H & a \\ C & a \\ (11) & bc \end{pmatrix}$.

Theorem 1.1([8]) There exist metrical N-linear connections $D\Gamma(N)$ on OscM, with respect to the homogeneous prolongation $\mathring{\mathbb{G}}$, which depend only on the metric

The Gauss-Weingarten formulae

 $g_{ab}(x,y)$. One of these connections has the "horizontal" coefficients

$$\begin{aligned}
\overset{H}{L}{}_{(00)}{}^{a}{}_{bc} &= \frac{1}{2}g^{ad} \left(\delta_{b}g_{dc} + \delta_{c}g_{bd} - \delta_{d}g_{bc}\right) \\
\overset{V}{L}{}_{(10)}{}^{a}{}_{bc} &= \frac{1}{2}h^{ad} \left(\delta_{b}h_{dc} + \delta_{c}h_{bd} - \delta_{d}h_{bc}\right)
\end{aligned}$$
(11)

and the "vertical" coefficients:

It is called the **Cartan metrical N-linear connection**. This linear connection will be used throughout this paper.

For this N-linear connection, we have the operators $\stackrel{H}{D}$ and $\stackrel{V}{D}$ which are given by the following relations

$$\begin{array}{l}
\overset{H}{D}X^{a} = dX^{a} + \overset{H}{\omega_{b}^{a}}X^{b} \\
\overset{V}{D}X^{a} = dX^{a} + \overset{V}{\omega_{b}^{a}}X^{b}.
\end{array}$$
(13)

We call these operators the **horizontal** and **vertical covariant differentials**. The 1-forms which define these operators will be called the **horizontal** and **vertical 1-form**, where

$$\begin{aligned}
\overset{H_{a}}{\omega_{b}} &= \overset{H}{\underset{(00)}{\overset{a}{bc}}} dx^{c} + \overset{H}{\underset{(01)}{\overset{a}{bc}}} \delta y^{c} \\
\overset{V}{\omega_{b}}^{a} &= \overset{V}{\underset{(10)}{\overset{a}{bc}}} dx^{c} + \overset{V}{\underset{(11)}{\overset{a}{bc}}} \delta y^{c}.
\end{aligned} \tag{14}$$

We have

Theorem 1.2[16] The d-tensors of torsion of the Cartan metrical N-linear connection D have the next expressions:

$$\begin{array}{l}
 H \\
 T \\
 (00) bc \\
 bc \\
 = \begin{pmatrix} H \\ 00 \end{pmatrix} bc \\
 = \begin{pmatrix} V \\ 01 \end{pmatrix} bc$$

bases $\{\delta_a, \dot{\partial}_{1a}\}$, the following d-tensors of curvature "horizontals"

and the "verticals"

3 The relative covariant derivatives

Let $D\Gamma(N)$, the Cartan metrical N-linear connection of the manifold OscM. A classical method to determine the laws of derivation on a Finsler submanifold is the type of the coupling.

Theorem 2.1 The coupling of the N-linear connection D to the induced nonlinear connection \check{N} along \widetilde{OscM} is locally given by the set of coefficients $\check{D}\Gamma(\check{N}) =$

$$\begin{pmatrix} H & V & H & V \\ \check{L}^{a}_{b\delta}, \check{L}^{a}_{b\delta}, \check{C}^{a}_{b\delta}, \check{C}^{a}_{(11)}_{b\delta} \end{pmatrix}, where$$

$$\begin{cases}
\overset{H}{\check{L}}\overset{a}{}_{\delta\delta} = \overset{H}{\underset{(00)}{}^{b}b} \overset{a}{B} \overset{d}{\delta} + \overset{H}{\underset{(01)}{}^{c}b} \overset{a}{B} \overset{d}{\delta} K^{\bar{\delta}} \\
\overset{V}{\check{L}}\overset{a}{}_{\delta} = \overset{V}{\underset{(10)}{}^{a}b} \overset{a}{B} \overset{d}{\delta} + \overset{V}{\underset{(11)}{}^{a}b} \overset{a}{B} \overset{d}{\delta} K^{\bar{\delta}} \\
\overset{H}{\check{C}}\overset{a}{}_{\delta} = \overset{H}{\underset{(01)}{}^{c}b} \overset{a}{B} \overset{d}{\delta} \\
\overset{V}{\check{C}}\overset{a}{}_{\delta} = \overset{V}{\underset{(11)}{}^{a}b} \overset{a}{B} \overset{d}{\delta} \\
\overset{V}{\underset{(11)}{}^{c}b} = \overset{V}{\underset{(11)}{}^{a}b} \overset{a}{B} \overset{d}{\delta}.
\end{cases} (18)$$

Definition 2.2 We call the **induced tangent connection** on OscM by the metrical N-linear connection D, the couple of operators D^{\top} , D^{\top} which are defined by

$$\begin{split} \stackrel{H}{D^{\top}} & X^{\alpha} = B^{\alpha}_{b} \stackrel{H}{D} X^{b}, \\ & for X^{a} = B^{a}_{\gamma} X^{\gamma} \\ D^{\top} & X^{\alpha} = B^{\alpha}_{b} \stackrel{V}{D} X^{b}, \end{split}$$

where

$$D^{\top} X^{\alpha} = dX^{\alpha} + X^{\beta} \overset{H_{\alpha}}{\omega_{\beta}}$$
$$V^{\top} X^{\alpha} = dX^{\alpha} + X^{\beta} \overset{V_{\alpha}}{\omega_{\beta}}$$

and $\overset{H_{\alpha}}{\omega_{\beta}}, \overset{V_{\alpha}}{\omega_{\beta}}$ are called the **tangent connection 1-forms**. We have

Theorem 2.3 The tangent connections 1-forms are as follows:

$$\begin{aligned}
\overset{H}{\omega}{}_{\beta}^{\alpha} &= \overset{H}{\underset{(00)}{}_{\beta\delta}} du^{\delta} + \overset{H}{\underset{(01)}{}_{\beta\delta}} \delta v^{\delta} \\
\overset{V}{\omega}{}_{\beta}^{\alpha} &= \overset{V}{\underset{(10)}{}_{\beta\delta}} du^{\delta} + \overset{V}{\underset{(11)}{}_{\beta\delta}} \delta v^{\delta},
\end{aligned}$$
(19)

where

$$\begin{array}{l}
\overset{H}{L}{}^{\alpha}{}_{(00)}{}_{\beta\delta} = B^{\alpha}_{d} \left(B^{d}_{\beta\delta} + B^{f}_{\beta} \overset{H}{\check{L}}{}^{d}_{(00)}{}^{f}_{\delta} \right), \\
\overset{V}{L}{}^{\alpha}{}_{(10)}{}_{\beta\delta} = B^{\alpha}_{d} \left(B^{d}_{\beta\delta} + B^{f}_{\beta} \overset{V}{\check{L}}{}^{d}_{(10)}{}^{f}_{\delta} \right), \\
\overset{H}{C}{}^{\alpha}{}_{(01)}{}_{\beta\delta} = B^{\alpha}_{d} B^{f}_{\beta} \overset{H}{\check{C}}{}^{d}_{(01)}{}^{f}_{\delta}, \\
\overset{V}{}^{C}{}^{\alpha}{}_{(11)}{}_{\beta\delta} = B^{\alpha}_{d} B^{f}_{\beta} \overset{V}{}^{\check{C}}{}^{d}_{(11)}{}^{f}_{\delta}.
\end{array}$$
(20)

Definition 2.4 We call the *induced normal connection* on \widetilde{OscM} by the metrical N-linear connection D, the couple of operators D^{\perp} , D^{\perp} which are defined by

$$\begin{array}{l} {}^{H}_{D^{\perp}}X^{\overline{\alpha}} = B^{\alpha}_{b}\check{D}X^{b} \\ for \ X^{a} = B^{\alpha}_{\bar{\gamma}}X^{\bar{\gamma}} \\ D^{\perp}X^{\overline{\alpha}} = B^{\alpha}_{b}\check{D}X^{b}, \end{array}$$

where

$$\begin{split} & \stackrel{H}{D^{\perp}} X^{\overline{\alpha}} = dX^{\overline{\alpha}} + X^{\overline{\beta}} \stackrel{H}{\omega}_{\overline{\beta}}^{\overline{\alpha}} \\ & \stackrel{V}{D^{\perp}} X^{\overline{\alpha}} = dX^{\overline{\alpha}} + X^{\overline{\beta}} \stackrel{V}{\omega}_{\overline{\beta}}^{\overline{\alpha}} \end{split}$$

and $\overset{H}{\omega}_{\overline{\beta}}^{\overline{\alpha}}$, $\overset{V}{\omega}_{\overline{\beta}}^{\overline{\alpha}}$ are called the **normal connection 1-forms**. We have

Theorem 2.5 The normal connections 1-forms are as follows:

$$\overset{H}{\omega} \frac{\overline{\alpha}}{\overline{\beta}} = \overset{H}{\underset{(00)}{\Sigma}} \frac{\overline{\alpha}}{\overline{\beta}\delta} du^{\delta} + \overset{H}{\underset{(01)}{C}} \frac{\overline{\alpha}}{\overline{\beta}\delta} \delta v^{\delta}$$

$$\overset{V}{\omega} \frac{\overline{\alpha}}{\overline{\beta}} = \overset{V}{\underset{(10)}{\Sigma}} \frac{\overline{\alpha}}{\overline{\beta}\delta} du^{\delta} + \overset{V}{\underset{(11)}{C}} \frac{\overline{\alpha}}{\overline{\beta}\delta} \delta v^{\delta},$$
(21)

where

Now, we can define the relative (or mixed) covariant derivatives $\stackrel{H}{\nabla}$ and $\stackrel{V}{\nabla}$. **Theorem 2.6** The relative covariant (mixed) derivatives in the algebra of mixed d-tensor fields are the operators $\stackrel{H}{\nabla}$, $\stackrel{V}{\nabla}$ for which the following properties hold:

$$\begin{array}{l} \overset{H}{\nabla} f = df, \\ & \forall f \in \mathcal{F}\left(\widetilde{Osc\check{M}}\right) \\ \overset{V}{\nabla} f = df, \end{array}$$

$$\begin{split} \overset{H}{\nabla} X^{a} &= \overset{H}{\check{D}} X^{a}, \quad \overset{H}{\nabla} X^{\alpha} = \overset{H}{D^{\intercal}} X^{\alpha}, \quad \overset{H}{\nabla} X^{\overline{\alpha}} = \overset{H}{D^{\bot}} X^{\overline{\alpha}}, \\ \overset{V}{\nabla} X^{a} &= \overset{V}{\check{D}} X^{a}, \quad \overset{V}{\nabla} X^{\alpha} = \overset{V}{D^{\intercal}} X^{\alpha}, \quad \overset{V}{\nabla} X^{\overline{\alpha}} = \overset{H}{D^{\bot}} X^{\overline{\alpha}}. \end{split}$$
$$\overset{H}{\check{\omega}} \overset{V}{}_{b}, \overset{H}{\check{\omega}} \overset{V}{}_{b}, \overset{W}{\omega} \overset{H}{}_{\beta}, \overset{W}{\omega} \overset{W}{}_{\overline{\alpha}}, \overset{W}{}_{\overline{\alpha}} \text{ are called the connection 1-forms of } \overset{H}{\nabla}, \overset{V}{\nabla}. \end{split}$$

4 The Gauss-Weingarten formulae

As usual in the theory of the submanifolds we are interesed in finding the moving equations of the moving frame \mathcal{R} along OscM.

These equations, called also Gauss-Weingarten formulae, are obtained when the relative covariant derivatives of the vector fields from \mathcal{R} are expressed again in the frame \mathcal{R} .

Thus we have

Theorem 3.1 The following Gauss-Weingarten formulae hold:

$$\nabla^{V_i} B^a_\alpha = B^a_{\bar{\delta}} \Pi^{V_i \bar{\delta}}_\alpha, \tag{23}$$

$$\nabla^{V_i} B^a_{\bar{\alpha}} = -B^a_\delta \Pi^{V_i}_{\bar{\alpha}},\tag{24}$$

where

$$\begin{aligned}
\stackrel{V_i}{\Pi}{}_{\alpha}^{\bar{\delta}} &= \stackrel{V_i}{\underset{(0)}{H}}{}_{\alpha}{}^{\bar{\delta}}{}_{\beta}du^{\beta} + \stackrel{V_i}{\underset{(1)}{H}}{}_{\alpha}{}^{\bar{\delta}}{}_{\beta}\deltav^{\beta} \\
\stackrel{V_i}{\Pi}{}_{\bar{\delta}}^{\alpha} &= g^{\alpha\sigma}\delta_{\bar{\delta}\bar{\sigma}}\stackrel{V_i}{\Pi}{}_{\sigma}^{\bar{\sigma}},
\end{aligned}$$
(25)

and the d-tensors

$$\begin{array}{l}
\overset{H}{H}_{(0)}{}_{\alpha}{}^{\bar{\delta}}{}_{\beta} = B_{d}^{\bar{\delta}}\left(B_{\alpha\beta}^{d} + B_{\alpha}^{f}\overset{H}{\check{L}}\overset{d}{}_{(00)}^{f}{}_{\beta}\right) & \overset{V}{H}_{\alpha}{}^{\bar{\delta}}{}_{\beta} = B_{d}^{\bar{\delta}}\left(B_{\alpha\beta}^{d} + B_{\alpha}^{f}\overset{V}{\check{L}}\overset{d}{}_{(10)}^{f}{}_{\beta}\right) \\
\overset{H}{H}_{\alpha}{}^{\bar{\delta}}{}_{\beta} = B_{d}^{\bar{\delta}}B_{\alpha}^{f}\overset{H}{\check{C}}\overset{d}{}_{\beta} & \overset{V}{H}_{\alpha}{}^{\bar{\delta}}{}_{\beta} = B_{d}^{\bar{\delta}}B_{\alpha}^{f}\overset{V}{\check{C}}\overset{d}{}_{\beta},
\end{array}$$
(26)

are the fundamental d-tensors of the second order of manifold \widetilde{OscM} , $(i = 0, 1, V_0 = H, V_1 = V)$. **Proof** From (11) and (12) we have

$$\begin{split} \overset{H}{\nabla} B^{a}_{\alpha} &= B^{a}_{\alpha|0\beta} du^{\beta} + B^{a}_{\alpha}|_{0\beta} \,\delta v^{\delta} \\ &= \left(\frac{\delta B^{a}_{\alpha}}{\delta u^{\beta}} + \overset{H}{\overset{L}{\overset{a}{\overset{a}{b}}}} B^{b}_{\alpha} - \overset{H}{\overset{L}{\overset{b}{(00)}}} \overset{\delta}{\overset{\alpha}{\overset{\beta}{\beta}}} B^{a}_{\delta} \right) du^{\beta} + \\ &+ \left(\frac{\delta B^{a}_{\alpha}}{\delta v^{\beta}} + \overset{H}{\overset{C}{\overset{a}{0}}} B^{b}_{\alpha} - \overset{H}{\overset{C}{(01)}} \overset{\delta}{\overset{\alpha}{\overset{\beta}{\beta}}} B^{a}_{\delta} \right) \delta v^{\beta} \\ &= B^{a}_{\alpha\beta} du^{\beta} + B^{b}_{\alpha} \left(\overset{H}{\overset{L}{\overset{a}{b}}} au^{\beta} + \overset{H}{\overset{C}{(01)}} \overset{a}{\overset{\beta}{\overset{\beta}{\beta}}} \delta v^{\beta} \right) - \\ &- B^{a}_{\delta} \left[B^{\delta}_{d} \left(B^{d}_{\alpha\beta} + B^{f}_{\alpha} \overset{H}{\overset{L}{\overset{\beta}{b}}} du^{\beta} + B^{\delta}_{d} B^{f}_{\alpha} \overset{H}{\overset{C}{(01)}} \overset{d}{\overset{\beta}{\overset{\beta}{\beta}}} \delta v^{\beta} \right) \right] \\ = D^{a}_{\delta} \left[B^{\delta}_{d} \left(B^{d}_{\alpha\beta} + B^{f}_{\alpha} \overset{H}{\overset{L}{\overset{\beta}{b}}} du^{\beta} + B^{\delta}_{d} B^{f}_{\alpha} \overset{H}{\overset{C}{(01)}} \overset{d}{\overset{\beta}{\overset{\beta}{\beta}}} \delta v^{\beta} \right) \right] \right] \\ = D^{a}_{\delta} \left[B^{\delta}_{d} \left(B^{d}_{\alpha\beta} + B^{f}_{\alpha} \overset{H}{\overset{L}{\overset{\beta}{b}}} du^{\beta} + B^{\delta}_{d} B^{f}_{\alpha} \overset{H}{\overset{C}{(01)}} dv^{\beta}} \right) du^{\beta} + B^{\delta}_{d} B^{f}_{\alpha} \overset{H}{\overset{C}{(01)}} dv^{\beta}} dv^{\beta} \right] \right] \\ = D^{a}_{\delta} \left[B^{\delta}_{d} \left(B^{d}_{\alpha\beta} + B^{f}_{\alpha} \overset{H}{\overset{L}{\overset{\beta}{\delta}} dv^{\beta} \right) du^{\beta} + B^{\delta}_{d} B^{f}_{\alpha} \overset{H}{\overset{C}{\delta}} dv^{\beta} dv^{\beta} \right] \right] \\ = D^{a}_{\delta} \left[B^{\delta}_{d} \left(B^{d}_{\alpha\beta} + B^{f}_{\alpha} \overset{H}{\overset{L}{\overset{\beta}{\delta}} dv^{\beta} \right) du^{\beta} + B^{\delta}_{d} B^{f}_{\alpha} \overset{H}{\overset{C}{\delta}} dv^{\beta} dv^{\beta} \right) \right] \\ = D^{a}_{\delta} \left[B^{\delta}_{d} \left(B^{d}_{\alpha\beta} + B^{f}_{\alpha} \overset{H}{\overset{L}{\overset{\beta}{\delta}} dv^{\beta} \right) du^{\beta} + B^{\delta}_{d} B^{\delta}_{\alpha} \overset{H}{\overset{\delta}{\delta}} dv^{\beta} dv^{\beta} \right) \right] \\ = D^{a}_{\delta} \left[B^{\delta}_{\delta} \left(B^{d}_{\alpha\beta} + B^{f}_{\alpha} \overset{H}{\overset{L}{\overset{\delta}{\delta}} dv^{\beta} \right) du^{\beta} + B^{\delta}_{\delta} B^{\delta}_{\alpha} \overset{H}{\overset{\delta}{\delta}} dv^{\beta} dv^{\beta} \right] \right] \\ = D^{a}_{\delta} \left[B^{\delta}_{\delta} \left(B^{d}_{\alpha\beta} + B^{\delta}_{\alpha\beta} \overset{H}{\overset{L}{\overset{\delta}{\delta}} dv^{\beta} \right) du^{\beta} + B^{\delta}_{\delta} B^{\delta}_{\alpha} \overset{H}{\overset{K}{\overset{\delta}{\delta}} dv^{\beta} dv^{\beta} \right] \right] \\ = D^{a}_{\delta} \left[B^{\delta}_{\delta} \left(B^{d}_{\alpha\beta} + B^{\delta}_{\alpha\beta} \overset{H}{\overset{K}{\overset{K}{\delta}} dv^{\beta} \right) \right] \\ = D^{a}_{\delta} \left[B^{\delta}_{\delta} \left(B^{\delta}_{\alpha\beta} + B^{\delta}_{\alpha\beta} \overset{H}{\overset{K}{\overset{K}{\delta}} dv^{\beta} \right) \right] \\ = D^{a}_{\delta} \left[B^{\delta}_{\alpha\beta} & B^{\delta}_{\alpha\beta} \overset{H}{\overset{K}{\delta}} \dot{v}^{\beta} \right] \\ = D^{a}_{\delta} \left[B^{\delta}_{\alpha\beta} & B^{\delta}_{\alpha\beta} \overset{H}{\overset{K}{\delta} v^{\beta} \right] \right] \\ = D^{a}_{\delta} \left[B^{\delta}_{\alpha\beta} & B^{\delta}_{\alpha\beta} \overset{H}{\overset{K}{\delta}$$

Using (25) we get relation (23) for $V_0 = H$.

Now, by applying $\stackrel{H}{\nabla}$ to both sides of the equations

$$g_{ab}B^a_\alpha B^b_{\bar{\beta}} = 0$$

one gets

$$g_{ab}B^a_{\bar{\delta}}{}^H_{\alpha}{}^{\bar{\delta}}B^b_{\bar{\beta}} + g_{ab}B^a_{\alpha}{}^H_{\alpha}B^b_{\bar{\beta}} = 0.$$

Multiplying these relation with B^{α}_d we obtain

$$g_{bd} \stackrel{H}{\nabla} B^b_{\bar{\beta}} - B^a_{\bar{\delta}} B^{\bar{\delta}}_d g_{ab} \stackrel{H}{\nabla} B^b_{\bar{\beta}} = -B^{\alpha}_d \delta_{\bar{\beta}\bar{\gamma}} \stackrel{H}{\Pi^{\bar{\gamma}}_{\alpha}}.$$

But $B_{\bar{\delta}}^{a} B_{d}^{\bar{\delta}} g_{ab} \nabla^{H} B_{\bar{\beta}}^{b} = 0$. Consequently, we obtain the relations (24) for $V_{0} = H$. Analogously, for the operator ∇^{V} one gets the other relations.

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