Bulletin of the *Transilvania* University of Braşov • Vol 9(58), No. 2 - 2016 Series III: Mathematics, Informatics, Physics, 25-40

DISTINGUISHED CURVATURES IN EXTENDED RELATIVISTIC DYNAMICS

Mircea NEAGU¹ and Alexandru OAN \breve{A}^2

Abstract

In this paper we present the distinguished (d-) curvatures for a Lagrangian inspired by relativistic optics in non-uniform media.

2010 Mathematics Subject Classification: 53C60, 53C80, 83C10. Key words: Minkowski metric, anisotropic optics, nonlinear connection, Cartan linear connection, d-torsions, d-curvatures.

1 Introduction

In geometrical optics [5], a special role is played by the Synge-Beil metric (see [1], [2], [4], [7]–[11])

$$g_{\alpha\beta}(x,y) = \varphi_{\alpha\beta}(x) + \gamma^2 y_{\alpha} y_{\beta}, \qquad (1)$$

where $\gamma(x) \geq 0$ is a positive smooth function on the space-time M^4 , and $\varphi_{\alpha\beta}(x)$ is a pseudo-Riemannian metric on M^4 . One assumes that the manifold M^4 (which is connected, simply connected and has dim $M^4 = 4$) is endowed with the local coordinates $(x^{\alpha})_{\alpha=\overline{1,4}} = (x^1 = t, x^2, x^3, x^4)$; for simplicity we use the system of units where the light velocity is c = 1. Obviously, the following rule holds: $y_{\alpha} = \varphi_{\alpha\mu}y^{\mu}$. Since the components of $\varphi_{\alpha\beta}(x)$ are dimensionless, the same are γy_{α} ; so we have $[\varphi_{\alpha\beta}(x)] = 1$, $[\gamma y_{\alpha}] = 1$.

In such a context, let us restrict our geometric-physical study to the Minkowski manifold $\mathcal{M}^4 = (\mathbb{R}^4, \eta_{ij})$ which has the local coordinates $(x) := (x^i)_{i=\overline{1,4}}$. It follows that the dimension of the corresponding tangent bundle $T\mathbb{R}^4$ is equal to eight, and its local coordinates are³

$$(x,y) := (x^i, y^i)_{i=\overline{1,4}} = (\underbrace{x^1, x^2, x^3, x^4}_{\text{space-time coordinates}}, \underbrace{y^1, y^2, y^3, y^4}_{\text{tangent vector}})$$

¹Department of Mathematics-Informatics, *Transilvania* University of Braşov, Romania, B-dul Iuliu Maniu, 50, e-mail: mircea.neagu@unitbv.ro

²Department of Mathematics-Informatics, *Transilvania* University of Braşov, Romania, B-dul Iuliu Maniu, 50, e-mail: alexandru.oana@unitbv.ro

³In this paper the Latin letters i, j, k, ... run from 1 to 4. The Einstein convention of summation is adopted all over this work.

Emerging from formula (1), we introduce the following metric on $T\mathbb{R}^4$, which is inspired by the J.L. Synge optics framework for the non-uniform medium:

$$\mathfrak{g}_{ij}(x,y) = \eta_{ij} + \gamma^2(x)y_iy_j, \qquad (2)$$

where $\eta = (\eta_{ij}) = \text{diag}(-1, 1, 1, 1)$ is the Minkowski metric, and $y_i = \eta_{ir}y^r$. Usually, we have $\gamma^2(x) = n^2(x) - 1$, where n = n(x) is the refractive index of the non-uniform medium (see [2], [7]–[10]). Using the metric (2), in what follows we will examine the special case of a possible anisotropic relativistic dynamical model (suggested in private discussions by physicist V.M. Red'kov), which is governed by the Lagrangian (in this model one considers that the particle has the mass m = 1) (see also [9])

$$L(x,y) = \frac{1}{2}\mathfrak{g}_{ij}(x,y)y^{i}y^{j}$$

$$= \frac{1}{2}(\eta_{ij} + \gamma^{2}y_{i}y_{j})y^{i}y^{j}$$

$$= \frac{1}{2}\eta_{ij}y^{i}y^{j} + \frac{\gamma^{2}}{2}||y||^{4},$$

$$(y^{2})^{2} + (y^{3})^{2} + (y^{4})^{2} = \eta_{ij}y^{i}y^{j}.$$
(3)

where $||y||^2 = -(y^1)^2 + (y^2)^2 + (y^3)^2 + (y^4)^2 = \eta_{ij}y^iy^j$.

Remark 1. Suppose that the refractive index n(x) is invariant with respect to Lorentz transformations. Since the Minkowski metric η_{ij} is invariant with respect to the linear transformations of coordinates induced by the Lorentz group O(3,1), it immediately follows that the Lagrangian (3) has a global geometrical character with respect to these Lorentz transformations.

Remark 2. A similar 3-dimensional anisotropic non-relativistic Lagrangian in which the Minkowski metric $(\eta_{ij})_{i,j=\overline{1,4}}$ is replaced with Euclidian metric $(\delta_{ij})_{i,j=\overline{1,3}}$ is studied by Neagu, Oana and Red'kov in paper [10]. That Lagrangian is invariant with respect to the orthogonal group O(3) and governs the **non-relativistic** extended dynamics.

Following the geometrical ideas from Lagrangian geometry of tangent bundles [6] or jet bundles [3], we further construct the pseudo-Riemann-Lagrange geometrical objects, such as the canonical nonlinear connection, the Cartan canonical linear connection, together with its d-torsions and d-curvatures, naturally associated with the Lagrangian (3).

2 Geometrical objects in relativistic extended dynamics

The Lagrangian (3) produces the fundamental metrical distinguished tensor

$$g_{ij}(x,y) = \frac{1}{2} \frac{\partial^2 L}{\partial y^i \partial y^j} = \sigma(x,y)\eta_{ij} + 2\gamma^2(x)y_i y_j,$$

where $\sigma(x, y) = (1/2) + \gamma^2(x)||y||^2$. Working on the domains of $T\mathbb{R}^4$ in which $\sigma(x, y) \neq 0$ and $\tau(x, y) = (1/2) + 3\gamma^2(x)||y||^2 \neq 0$, then the inverse matrix $[g^{-1}] = (g^{jk})_{i,k=\overline{1,4}}$ has the components

$$g^{jk}(x,y) = \frac{1}{\sigma(x,y)} \eta^{jk} - \frac{2\gamma^2(x)}{\sigma(x,y) \cdot \tau(x,y)} y^j y^k,$$

where $\eta^{jk} = \eta_{jk}$.

Following the Miron and Anastasiei geometrical ideas from the book [6], we deduce that, for the anisotropic Lagrangian (3), the associated *canonical nonlinear connection* $N = \left(N_j^i\right)$ on the tangent bundle $T\mathbb{R}^4$ has the local components (these are computed in [9])

$$N_{j}^{i} = \frac{2\gamma}{\sigma} y^{i} y_{j} \gamma_{0} + \frac{\gamma ||y||^{2}}{\sigma} \left(\delta_{j}^{i} \gamma_{0} + y^{i} \gamma_{j} - \gamma^{i} y_{j} - \frac{2\gamma^{2}}{\sigma} y^{i} y_{j} \gamma_{0} - \frac{6\gamma^{2}}{\tau} y^{i} y_{j} \gamma_{0} \right)$$

$$+ \frac{\gamma^{3} ||y||^{4}}{2\sigma} \left[\frac{1}{\sigma} \gamma^{i} y_{j} - \frac{3}{\tau} y^{i} \gamma_{j} - \frac{3}{\tau} \delta_{j}^{i} \gamma_{0} + \frac{6\gamma^{2}}{\sigma\tau^{2}} (\tau + 3\sigma) y^{i} y_{j} \gamma_{0} \right],$$

$$(4)$$

where $\gamma_s = \partial \gamma / \partial x^s$, $\gamma_{rs} = \partial^2 \gamma / \partial x^r \partial x^s$ and $\gamma^i = \eta^{ir} \gamma_r$, $\gamma_0 = \gamma_r y^r$.

Remark 3. In a uniform medium with the constant refractive index $n(x) = n \in [1, \infty)$, we have $\gamma_s = 0$. Consequently, in this case we have $N_i^i = 0$.

The nonlinear connection (4) produces the dual *adapted bases* of *d*-vector fields

$$\left\{\frac{\delta}{\delta x^{i}} = \frac{\partial}{\partial x^{i}} - N_{i}^{r} \frac{\partial}{\partial y^{r}} ; \frac{\partial}{\partial y^{i}}\right\} \subset \mathfrak{X}(T\mathbb{R}^{4})$$
(5)

and of d-covector fields

$$\left\{ dx^i \; ; \; \delta y^i = dy^i + N^i_r dx^r \right\} \subset \mathfrak{X}^*(T\mathbb{R}^4).$$
(6)

The description of all geometrical objects on the tangent space $T\mathbb{R}^4$ (e.g., the Cartan canonical linear connection and its torsion and curvature) will be made in local adapted components, with respect to the adapted bases (5) and (6).

For instance, by using the derivative operators (5) and the notations $N_{ij} := N_i^r \eta_{rj}$, $N_0^i = N_r^i y^r$, $N_{i0} := N_{ir} y^r$, $N_{0j} := N_{rj} y^r$, $N_{00} := N_{ij} y^i y^j$, by direct local computations, we find the adapted local components of the Cartan canonical connection $C\Gamma(N) = (L_{jk}^i, C_{jk}^i)$ associated with the Lagrangian (3), which are computed in [9]:

$$L_{jk}^{i} = -\frac{\gamma}{\sigma} \left[\gamma \left(\delta_{j}^{i} N_{k0} + \delta_{k}^{i} N_{j0} - \eta^{ir} \eta_{jk} N_{r0} \right) + ||y||^{2} \left(\eta_{jk} \gamma^{i} - \delta_{j}^{i} \gamma_{k} - \delta_{k}^{i} \gamma_{j} \right) + + \gamma \left\{ \left(N_{jk} + N_{kj} \right) y^{i} + \left(N_{k}^{i} - \eta^{ir} N_{rk} \right) y_{j} + \left(N_{j}^{i} - \eta^{ir} N_{rj} \right) y_{k} \right\} + + 2 \left(\gamma^{i} y_{j} y_{k} - y^{i} y_{j} \gamma_{k} - y^{i} y_{k} \gamma_{j} \right) \right] + \frac{2\gamma^{3} y^{i}}{\sigma \tau} \left[\gamma \left(y_{j} N_{k0} + y_{k} N_{j0} - \eta_{jk} N_{00} \right) + + ||y||^{2} \left(\delta_{jk} \gamma_{r} y^{r} - y_{j} \gamma_{k} - y_{k} \gamma_{j} \right) + 2 \left(y_{j} y_{k} \gamma_{r} y^{r} - y_{j} \gamma_{k} ||y||^{2} - y_{k} \gamma_{j} ||y||^{2} \right) + + \gamma \left\{ \left(N_{jk} + N_{kj} \right) ||y||^{2} + \left(N_{k0} - N_{0k} \right) y_{j} + \left(N_{j0} - N_{0j} \right) y_{k} \right\} \right],$$
(7)

$$C_{jk}^{i} = \frac{\gamma^{2}}{\sigma} \left(y^{i} \eta_{jk} + \delta_{j}^{i} y_{k} + \delta_{k}^{i} y_{j} \right) - \frac{2\gamma^{4}}{\sigma \tau} \left(||y||^{2} \eta_{jk} + 2y_{j} y_{k} \right) y^{i}.$$

$$\tag{8}$$

Moreover, the local components of the torsion tensor of the Cartan canonical N-linear connection produced by the anisotropic optical Lagrangian (3) are given by R_{jk}^i , P_{jk}^i , C_{jk}^i , where (these are also computed in [9])

$$\begin{split} R_{jk}^{i} &= y^{i} \left(y_{j} \gamma_{0k} - y_{k} \gamma_{0j} \right) \varphi + \left(\delta_{j}^{i} \gamma_{0k} - \delta_{k}^{i} \gamma_{0j} \right) \varepsilon + \eta^{is} (\gamma_{sk} y_{j} - \gamma_{sj} y_{k}) \omega \\ &+ y^{i} (y_{j} \gamma_{k} - \gamma_{j} y_{k}) (\gamma_{s} \gamma^{s}) [\frac{7\varphi}{\gamma} - \frac{6}{\tau^{2}} - 18\varphi\varepsilon + \varphi^{2} \|y\|^{2} + 2\varphi\omega - \varepsilon\phi \|y\|^{2}] \\ &+ (\delta_{j}^{i} \gamma_{k} - \delta_{k}^{i} \gamma_{j}) (\gamma_{s} \gamma^{s}) [\frac{2\varepsilon}{\gamma} - \frac{\|y\|^{2}}{\sigma\tau} - 9\varepsilon^{2} - \varphi\varepsilon \|y\|^{2}] \\ &+ \gamma^{i} (y_{j} \gamma_{k} - \gamma_{j} y_{k}) [\frac{\omega}{\gamma} + \frac{2\|y\|^{2}}{\sigma^{2}} + \frac{2\omega}{\sigma\tau} + \frac{\varepsilon\gamma\|y\|^{2}}{2\sigma^{3}} + \omega^{2}] \\ &- (\delta_{j}^{i} y_{k} - \delta_{k}^{i} y_{j}) (\gamma_{s} \gamma^{s})^{2} \varphi(\varphi \|y\|^{2} + \omega + \varepsilon) - (\delta_{j}^{i} y_{k} - \delta_{k}^{i} y_{j}) (\gamma_{s} \gamma^{s}) \omega\varepsilon, \end{split}$$

$$\begin{split} P_{jk}^{i} &= y^{i}y_{j}y_{k}\gamma_{0}[12\varphi\gamma^{2}\tau(\frac{1}{\sigma}-\sigma^{2})-\frac{9\gamma^{3}}{\sigma^{2}\tau}-\frac{4\gamma^{3}}{\sigma\tau}]+(\delta_{j}^{i}y_{k}+\delta_{k}^{i}y_{j})\gamma_{0}\varphi \\ &+y^{i}(y_{j}\gamma_{k}+\gamma_{j}y_{k})[\varphi-\frac{2\gamma}{\sigma}+\frac{6\gamma^{3}||y||^{2}}{\sigma\tau}]+y^{i}\eta_{jk}\gamma_{0}[\varphi-\frac{2\gamma^{3}||y||^{2}}{\sigma\tau}] \\ &+(\delta_{j}^{i}\gamma_{k}+\delta_{k}^{i}\gamma_{j})(-\frac{3\gamma^{3}||y||^{4}}{2\sigma\tau})+\gamma^{i}\eta_{jk}\frac{\gamma^{3}||y||^{4}}{2\sigma^{2}}+\gamma^{i}y_{j}y_{k}\frac{\gamma(4\sigma^{2}-1)}{2\sigma^{3}} \\ &+(\delta_{j}^{i}N_{k0}+\delta_{k}^{i}N_{j0}-\eta^{ii}\eta_{jk}N_{i0})\frac{\gamma^{2}}{\sigma}-y^{i}(N_{kj}+N_{jk})(\frac{2\gamma^{4}}{\sigma\tau}-\frac{\gamma^{2}}{\sigma}) \\ &+[(N_{j}^{i}-\eta^{ii}N_{ij})y_{k}+((N_{k}^{i}-\eta^{ii}N_{ik})y_{j})]\frac{\gamma^{2}}{\sigma} \\ &-(y^{i}N_{j0}y_{k}+y^{i}y_{j}N_{k0})\frac{4\gamma^{4}}{\sigma\tau}+(y^{i}\eta_{jk}N_{00}+y^{i}y_{j}N_{0k}+y^{i}N_{0j}y_{k})\frac{2\gamma^{4}}{\sigma\tau} \end{split}$$

where

$$\begin{split} \alpha &= 12\varphi\gamma^2\tau \left|\left|y\right|\right|^2 \left(\frac{1}{\sigma} - \sigma^2\right) - \frac{9\gamma^3\left|\left|y\right|\right|^2}{\sigma^2\tau} + 2\varphi - \frac{\gamma}{2\sigma^3},\\ \varepsilon &= \frac{\gamma\left|\left|y\right|\right|^2 (2\tau+1)}{4\sigma\tau}, \varphi = \frac{\gamma(12\sigma^2 - 6\sigma + 1)}{4\sigma^2\tau^2},\\ \phi &= \frac{\gamma^3(12\sigma - 3)}{\sigma^3\tau^2} - \frac{12\sigma^2\gamma^2\varphi}{\tau}, \omega = -\frac{\gamma\left|\left|y\right|\right|^2 (2\sigma + 1)}{4\sigma^2}. \end{split}$$

In the sequel, note that the local components of the curvature tensor of a general Cartan canonical N-linear connection, are given by the general formulas

(see Miron-Anastasiei's book [6])

$$\begin{split} R^{i}_{jkl} &= \frac{\delta L^{i}_{jk}}{\delta x^{l}} - \frac{\delta L^{i}_{jl}}{\delta x^{k}} + L^{r}_{jk}L^{i}_{rl} - L^{r}_{jl}L^{i}_{rk} + C^{i}_{jr}R^{r}_{kl} \\ P^{i}_{jkl} &= \frac{\partial L^{i}_{jk}}{\partial y^{l}} - C^{i}_{jl|k} + C^{i}_{jr}P^{r}_{kl}, \\ S^{i}_{jkl} &= \frac{\partial C^{i}_{jk}}{\partial y^{l}} - \frac{\partial C^{i}_{jl}}{\partial y^{k}} + C^{r}_{jk}C^{i}_{rl} - C^{r}_{jl}C^{i}_{rk}, \end{split}$$

where

$$C^i_{jl|k} \stackrel{def}{=} \frac{\delta C^i_{jl}}{\delta x^k} + C^r_{jl} L^i_{rk} - C^i_{rl} L^r_{jk} - C^i_{jr} L^r_{lk}.$$

Consequently, as a novelty of this paper, using the formulas (4), (7), (8) and the derivative operators (5), after very laborious and complicated calculations, we get that the Cartan canonical N-linear connection produced by the anisotropic optical Lagrangian (3) is characterized by *three* effective local curvature d-tensors, namely R_{jkl}^{i} , P_{jkl}^{i} , S_{jkl}^{i} , where

$$\begin{split} R_{jkl}^{i} &= y^{i} y_{j} \left(y_{k} \gamma_{l} - \gamma_{k} y_{l}\right) \left\{ (\gamma_{0})^{2} \left[\left(\omega + \varepsilon + \varphi \left\|y\right\|^{2} \right) \mathfrak{T}_{2} - \varepsilon \mathfrak{T}_{1} \right] + N_{00} \frac{4\gamma^{5} (4\gamma^{3} - 1)}{\sigma^{2} \tau^{2}} \right\} \\ &+ y^{i} y_{j} \left(\gamma_{k} y_{l} - y_{k} \gamma_{l} \right) \gamma_{0} \left[2 \left(\varphi - 6\gamma^{2} \right) \frac{\gamma}{\sigma \tau} + \left(\omega - 2\varepsilon \right) \mathfrak{T}_{3} + \frac{\gamma^{2}}{\sigma} \left(\frac{4\gamma}{\tau} + \varphi \right) \left(\omega - 3\varepsilon \right) - \\ &- \mathcal{A} - 16\varphi\gamma^{2} - \frac{8\varphi}{\sigma \tau^{2}} - \frac{4\gamma^{4}}{\sigma \tau} \left(\mathcal{A} \left\|y\right\|^{2} + \mathcal{B} + 2\mathcal{C} \right) + \mathcal{R}_{1} \frac{\gamma^{2}(\tau - \sigma)}{\sigma \tau} + \varphi \mathfrak{T}_{7} \\ &+ \left(\mathcal{R}_{1} \left\|y\right\|^{2} + \mathcal{R}_{2} + \mathcal{R}_{3} \right) \frac{4\gamma^{4}}{\sigma \tau} \right] \\ &+ y^{i} y_{j} \left(y_{k} \gamma_{0l} - \gamma_{0k} y_{l} \right) \left[\varphi \frac{\sigma \tau + \gamma^{2} (\sigma + \tau)}{\sigma \tau} + \frac{4\gamma^{3}}{\sigma \tau} \left(1 + \omega \right) - \frac{2\gamma^{2} \sigma (2\tau + 1)}{\sigma^{2} \tau^{2}} \right] \\ &+ y^{i} y_{j} \left(N_{0} N_{0l} - N_{0k} N_{0l} \right) \frac{8\gamma^{6} (\gamma^{2} - \tau)}{\sigma^{2} \tau^{2}} \\ &+ y^{i} y_{j} \left(N_{0k} y_{l} - y_{k} N_{0l} \right) \gamma_{0} \left[\left(\varphi \left\|y\right\|^{2} + \omega + \varepsilon \right) \frac{4\gamma^{6} (2\tau + 1)}{\sigma^{2} \tau^{2}} - \varphi \frac{6\gamma^{4}}{\sigma \tau} \\ &+ \frac{4\gamma^{5} (\sigma \tau - \|y\|^{2} + 2\gamma^{2} - \tau)}{\sigma^{2} \tau^{2}} \right] \\ &+ y^{i} y_{j} \left(N_{0k} \gamma_{l} - \gamma_{k} N_{0l} \right) \left[8\varphi\gamma^{3} - \frac{4\gamma^{2}}{\sigma \tau^{2}} - \frac{8\gamma^{3}}{\sigma \tau} + \varepsilon \frac{6\gamma^{4}}{\sigma \tau} - \varepsilon \left\|y\right\|^{2} \frac{4\gamma^{6} (2\tau + 1)}{\sigma^{2} \tau^{2}} + \frac{2\gamma^{5}}{\sigma^{2} \tau^{2}} \right] \\ &- \frac{\gamma^{3} (1 - \sigma)}{\sigma^{2} \tau^{2}} - \frac{\gamma^{3} [(4\sigma - 1)\tau - 2\gamma^{2} + 16\gamma^{5} \|y\|^{2}]}{\sigma^{2} \tau^{2}} \right\} - y^{i} y_{j} \left(N_{k0} y_{l} - y_{k} N_{l0} \right) \gamma_{0} \left\{ \left(\varphi \left\|y\right\|^{2} + \omega + \varepsilon \right) \frac{8\gamma^{6} (2\tau + 1)}{\sigma^{2} \tau^{2}} - \varphi \frac{12\gamma^{4}}{\sigma \tau} \\ &+ \frac{\gamma^{5} [(4 + 27\gamma^{2}) \|y\|^{2} + 8\sigma]}{\sigma^{2} \tau^{2}} - \frac{4\gamma^{5}}{\sigma^{2} \tau^{2}} \right\} \right\} \\ &+ y^{i} y_{j} \left(N_{k0} y_{l} - y_{k} N_{l0} \right) \gamma_{0} \left\{ \left(\varphi \left\|y\right\|^{2} + \omega + \varepsilon \right) \frac{8\gamma^{6} (2\tau + 1)}{\sigma^{2} \tau^{2}} - \varphi \frac{12\gamma^{4}}{\sigma \tau} \\ &+ \frac{\gamma^{5} [(4 + 27\gamma^{2}) \|y\|^{2} + 8\sigma]}{\sigma^{2} \tau^{2} \tau^{2}} - \frac{4\gamma^{5}}{\sigma^{2} \tau^{2}} \right\} \right\} \\ &+ y^{i} y_{j} \left(N_{rk} \gamma^{r} y_{l} - y_{k} N_{rl} \gamma^{r} \right) \frac{2\gamma^{3} (2\gamma^{2} - \tau)}{\sigma^{2} \tau^{2}} + y^{i} y_{j} \left(N_{rk} \gamma^{r} y_{l} - y_{k} N_{rl} \gamma^{r} \right) \frac{\gamma^{2} (2\gamma^{2} - \tau)}{\sigma^{2} \tau^{2}} + y^{i} y_{j} \left(N_{rk} \gamma^{r} y_{l} - \gamma^{k} N_{rl} \gamma^{r} \right) \frac{\gamma^{4} (2\gamma^{2} - \tau)}{\sigma^{2} \tau^{2}} + y^{i} y_{j} \left(N_{rk} \gamma^{r} y_{l}$$

$$\begin{split} & +y^{\dagger}y_{j}\left(N_{k}^{r}y_{l}-y_{k}N_{l}^{r}\right)\left(\gamma^{r}\frac{s_{j}^{2}}{\sigma^{2}\tau}+N_{r}6\frac{4\gamma^{2}}{\sigma^{2}\tau}\right) \\ & +y^{\dagger}y_{j}\left(N_{kl}-N_{lk}\right)\left[e^{\frac{2\gamma^{4}\left(1-|y||^{2}\right)}{\sigma^{2}}-\frac{18\gamma^{2}||y||^{2}}{\sigma\tau}-24\varphi\gamma^{2}||y||^{2}+\frac{12\gamma||y||^{2}}{\sigma\tau^{2}} \\ & +z\frac{\gamma^{2}(\varepsilon-\omega)}{\sigma\tau}+\Re_{3}\frac{\gamma^{2}}{\tau}+\frac{2(1-|y||^{2})}{(\sigma^{2}}-\frac{\gamma^{2}||y||^{2}}{\sigma^{2}\tau^{2}}+\Upsilon_{7}\left(2\omega-\varepsilon\right)-\Im_{3}\varepsilon||y||^{2}-\frac{2\gamma(\varepsilon-\omega)}{\sigma\tau}\right] \\ & +y^{\dagger}\gamma_{j}\left(N_{0k}y_{l}-y_{k}N_{0l}\right)\left[\frac{2\gamma^{4}\left(1-|y||^{2}\right)}{\sigma\tau}+\frac{\gamma^{2}}{\sigma^{2}\tau^{2}}+\frac{\gamma^{2}}{\sigma^{2}\tau^{2}}-\frac{1}{\sigma^{2}\tau^{2}}\right] \\ & +\left(y^{\dagger}\gamma_{jl}y_{k}-\gamma_{jk}y_{l}\right)\left[\frac{2\gamma^{4}\left(1-|y||^{2}\right)}{\sigma\tau}+\frac{\gamma^{2}}{\sigma^{2}\tau^{2}}+(\gamma^{2}-\sigma^{2})^{2}+16\gamma^{5}||y||^{2}}\right] \\ & -y^{\dagger}\gamma_{j}\left(N_{k0}y_{l}-y_{k}N_{0l}\right)\left[\frac{2\gamma^{4}\left(1-\frac{1}{\sigma\tau}\right)^{2}}{\sigma^{2}\tau}+\frac{\gamma^{3}(\tau-3-2-\gamma^{2}+16\gamma^{5}||y||^{2}}{\sigma^{2}\tau^{2}}\right] \\ & +y^{i}N_{j0}\left(y_{k}\eta_{l}-\gamma_{k}y_{l}\right)\left[\frac{16\gamma^{2}}{\sigma\tau}-16\varphi\gamma^{3}+\frac{8\gamma^{2}}{\sigma\tau^{2}}-\varepsilon^{\frac{12\gamma^{4}}{2}}+\varepsilon^{\frac{8\gamma^{6}}{2}||y||^{2}(2\tau+1)}\gamma_{0}\right. \\ & +\frac{\omega^{2}\gamma^{4}}{\sigma^{4}\tau}+\frac{\gamma^{3}(\tau-2)}{\sigma^{2}\tau^{2}}\right] \\ & +y^{i}N_{j0}\left(N_{k0}y_{l}-y_{k}N_{l0}\right)\frac{8\gamma^{6}\left(\tau-\gamma^{2}\right)}{\sigma^{2}\tau^{2}}-y^{i}N_{j0}\left(N_{0k}y_{l}-y_{k}N_{0l}\right)\frac{2\gamma^{6}\left(4\gamma^{2}-7\sigma+2\right)}{\sigma^{2}\tau^{2}}} \\ & +y^{i}N_{0j}\left(y_{k}\eta_{l}-\gamma_{k}y_{l}\right)\left[8\varphi\gamma^{3}+\frac{\gamma^{2}\left(\gamma-4\right)}{\sigma^{2}\tau^{2}}-\frac{8\gamma^{3}}{\sigma^{3}}+\frac{2\left(3\epsilon+\omega\right)\gamma^{*}}{\sigma\tau}} \\ & -\varepsilon\left(1y\right)\left[\frac{4\gamma^{6}\left(2\gamma+1\right)}{\sigma^{2}\tau^{2}}+\frac{2\gamma^{5}}{\sigma^{2}\tau^{2}}\right] \\ & +y^{i}\left(N_{jk}y_{l}-y_{k}\eta_{jl}\right)\gamma_{0}\left[\left(y||y||^{2}+\omega+\varepsilon\right)\left(\gamma^{5}+\varepsilon\frac{4\gamma^{4}}{\sigma\tau} \\ & -\varepsilon\left(1y\right)\left[\frac{2\gamma^{6}\left(2\gamma^{2}-\tau\right)}{\sigma^{2}\tau^{2}}\right] \\ & +y^{i}\left(\eta_{jk}\gamma_{l}-\gamma_{k}\eta_{jl}\right)\gamma_{0}\left[\frac{8}{\sigma}+\frac{6\gamma^{2}||y||^{2}}{\sigma\tau} \\ & +\varepsilon\left(1y\right)\left[\frac{2\gamma^{2}}{\sigma\tau^{2}}+\frac{2\gamma^{2}(2\gamma^{2}-\tau)}{\sigma\tau^{2}}\right] \\ & +y^{i}\left(\eta_{jk}\gamma_{0}-\gamma_{0k}\eta_{jl}\right)\left[\frac{\gamma^{2}}{\tau}\left(\omega+2\varepsilon+\varphi\left||y||^{2}\right)+\frac{2\gamma^{3}||y||^{2}}{\sigma\tau^{2}}+\varepsilon^{2}\left(2\gamma^{2}-\gamma\tau\right)}{\sigma\tau}\right] \\ & +y^{i}\left(\eta_{jk}\gamma_{0}-\gamma_{k}\eta_{jl}\right)N_{00}\left[\left(\omega+2\varepsilon+\varphi\left||y||^{2}\right)+\frac{2\gamma^{3}||y||^{2}}{\sigma\tau^{2}}+\varepsilon^{2}\left(2\gamma^{2}-\tau\right)}{\sigma\tau}\right] \\ & +y^{i}\left(\eta_{jk}y_{0}-\gamma_{k}\eta_{jl}\right)N_{00}\left[\left(\omega^{2}+\varepsilon+\varphi\left||y||^{2}\right)+\frac{2\gamma^{3}||y||^{2}}{\sigma\tau^{2}}+\varepsilon^{2}\left(2\gamma^{2}-\tau\right)}{\sigma\tau}\right] \\ & +y^{i}\left(\eta_{jk}y_{0}-\gamma_{k}\eta_{jl}\right)N_{00}\left(\left(\omega^{6}+\varepsilon+\varphi\left||y||^{2}\right)+\frac{2\gamma^{3}||y||^{2}}{\sigma\tau^{2}}}\right) \\ &$$

$$\begin{split} &+y^{i}\left(\eta_{jk}y_{ll}-y_{k}\eta_{jl}\right)\left[\eta^{rr}\left(N_{r}0\right)^{2}\frac{4\gamma^{5}}{\sigma^{2}\tau}-\eta^{rr}N_{r}0N_{0}r\frac{2\gamma^{5}}{\sigma^{2}\tau^{2}}\right] \\ &+y^{i}\left(\eta_{jk}N_{0l}-N_{0k}\eta_{jl}\right)\gamma_{0}\left(\frac{2\gamma^{4}}{\sigma^{4}\tau^{2}}+\frac{4\gamma^{2}||y||^{2}}{\sigma^{2}\tau^{2}}\right) \\ &+y^{i}\left(\eta_{jk}N_{0l}-N_{0k}\eta_{jl}\right)N_{00}\left(\frac{4\gamma^{6}}{\sigma^{2}\tau^{2}}+\frac{8\sigma^{7}||y||^{2}}{\sigma^{2}\tau^{2}}\right) \\ &+y^{i}\left(\eta_{jk}N_{0l}-N_{k0}\eta_{jl}\right)N_{00}\left(\frac{4\gamma^{6}(\sigma^{2}+2\tau^{2}-\gamma^{2})}{\sigma^{2}\tau^{2}}\right) \\ &+y^{i}\left(\eta_{jk}N_{0l}-N_{k0}\eta_{jl}\right)N_{00}\left(\frac{4\gamma^{6}(\sigma^{2}+2\tau^{2}-\gamma^{2})}{\sigma^{2}}\right) \\ &+y^{i}\left(\eta_{jk}N_{0l}-N_{k0}\eta_{jl}\right)N_{00}\left(\frac{4\gamma^{6}(\sigma^{2}+2\tau^{2}-\gamma^{2})}{\sigma^{2}}\right) \\ &+y^{i}\left(\eta_{jk}(N_{lr}+N_{rl})-\eta_{jl}\left(N_{rk}+N_{kr}\right)\right)\left[\eta^{rr}N_{0}\frac{\gamma^{4}\left(2\gamma^{2}-\tau\right)\left(2\gamma^{2}-\sigma^{2}-\tau\right)}{\sigma^{2}\tau^{2}}-\gamma^{2}\gamma^{3}\left(\frac{1}{\sigma^{2}\tau^{2}}\right)\right] \\ &-y^{i}\left(\delta_{jk}N_{0l}-\delta_{jl}N_{0k}\right)\gamma_{0}\frac{2\gamma^{4}}{\sigma^{2}}+y^{i}\left(\delta_{jk}N_{0l}-\delta_{jl}N_{k0}\right)\gamma_{0}\frac{4\gamma^{4}}{\sigma^{4}} \\ &+y^{i}\left[\left(N_{jk}+N_{kj}\right)N_{0l}-\left(N_{jl}+N_{lj}\right)N_{k0}\right]\frac{\gamma^{4}\left(2\gamma^{2}-\tau\right)\left(2\gamma^{2}-\sigma^{2}-\sigma^{2}\right)}{\sigma^{2}\tau^{2}} \\ &+y^{i}\left[\left(N_{jk}+N_{kj}\right)N_{0l}-\left(N_{jl}+N_{lj}\right)N_{k0}\right]\frac{\gamma^{4}\left(2\gamma^{2}-\tau\right)\left(2\gamma^{2}-\sigma^{2}-\sigma^{2}\right)}{\sigma^{2}\tau^{2}} \\ &+N_{j}^{i}\left(y_{k}\gamma_{l}-\gamma_{ky}y_{l}\right)\left[\left(\varepsilon-\omega\right)\frac{\gamma^{2}}{\sigma}-\frac{2\gamma^{4}\left[|y||^{2}}{\sigma^{2}}\right) \\ &+N_{j}^{i}\left(y_{k}\gamma_{l}-\gamma_{ky}y_{l}\right)\left[\left(\varepsilon-\omega\right)\frac{\gamma^{2}}{\sigma}-\frac{2\gamma^{4}\left[|y||^{2}}{\sigma^{2}}\right) \\ &+\left(N_{k}^{i}y_{j}y_{l}-N_{l}^{i}y_{j}y_{k}\right)\left(\frac{\gamma^{2}\left[2\sigma^{2}-\tau\right]}{\sigma^{2}}+\frac{2\gamma^{2}\left[|y||^{2}}{\sigma^{2}}\right] \\ &+\left(N_{k}^{i}y_{j}y_{l}-N_{l}^{i}y_{j}y_{k}\right)\left(\frac{\gamma^{2}\left[2\sigma^{2}+\gamma^{4}\left(\sigma^{2}-1\right)\right]}{\sigma^{2}}\right] \\ &+\left(N_{k}^{i}y_{j}y_{l}-N_{l}^{i}y_{j}y_{k}\right)\left(\frac{\gamma^{2}}{\sigma^{2}}+\frac{\gamma^{4}\left(\sigma^{2}-1\right)}{\sigma^{2}\tau^{2}}\right) \\ &+\left(N_{k}^{i}y_{j}y_{l}-N_{l}^{i}y_{j}y_{k}\right)\left(\frac{\gamma^{2}}{\sigma^{2}}+\frac{\gamma^{4}\left(\sigma^{2}-1\right)}{\sigma^{2}\tau^{2}}\right) \\ &+\left(N_{k}^{i}y_{j}y_{l}-N_{l}^{i}y_{k}y_{l}\right)\left(\frac{\gamma^{2}}{\sigma^{2}}+\frac{\gamma^{4}\left(\sigma^{2}-1\right)}{\sigma^{2}\tau^{2}}}\right) \\ &+\left(N_{k}^{i}y_{j}y_{l}-N_{l}^{i}y_{k}y_{l}\right)\left(\frac{\gamma^{2}}{\sigma^{2}}+\frac{\gamma^{4}\left(\sigma^{2}-1\right)}{\sigma^{2}\tau^{2}}}\right) \\ &+\left(N_{k}^{i}y_{j}y_{l}-N_{l}^{i}y_{k}y_{l}\right)\left(\frac{\gamma^{2}}{\sigma^{2}}+\frac{\gamma^{4}\left(\sigma^{2}-1\right)}{\sigma^{2}\tau^{2}}}\right) \\ &+\left(N_{k}^{i}y_{j}y_{l}-N_{l}^{i}y_{k}y_{l}\right)\left(\frac{\gamma^{2}}{\sigma^{2}}+\frac{\gamma^{4}\left(\sigma^{2}-1\right)}{\sigma^$$

$$\begin{split} &+\gamma^{i}\left(\eta_{jk}\gamma_{l}-\eta_{jl}\gamma_{k}\right)\left[\frac{(\sigma-1)||\mathbf{u}||^{2}}{\sigma^{2}}+\frac{\gamma^{2}||\mathbf{u}||^{2}}{\sigma^{2}}-\frac{(\gamma^{3}||\mathbf{u}||^{2}(3\gamma^{2}-\tau))}{\sigma^{2}}\right] \\ &-\gamma^{i}\left[(N_{jk}+N_{kj})\,y_{l}-(N_{jl}+N_{lj})\,y_{k}\right]\left[\omega^{\frac{\gamma^{2}(2\gamma^{2}-\tau)}{\sigma\tau}}-\frac{(\gamma^{3}||\mathbf{u}||^{2}(3\gamma^{2}-\tau))}{\sigma^{2}\tau}\right] \\ &+\eta^{ii}N_{i0}\left(\eta_{jk}y_{l}-\eta_{ij}\gamma_{k}\right)\left[\frac{\gamma^{2}||\mathbf{u}||^{2}(\gamma-2)}{\sigma^{2}}+\frac{\gamma}{\sigma^{2}}+\frac{\varepsilon\gamma^{2}(3\sigma-1)}{\sigma^{2}}\right] \\ &+\eta^{ii}N_{i0}\left(\eta_{jk}N_{l0}-\eta_{jl}N_{k0}\right)\frac{\gamma^{2}}{\sigma^{2}}+\eta^{ii}N_{i0}y_{j}\left(N_{k0}y_{l}-y_{k}N_{l0}\right)\frac{4\gamma}{\sigma^{2}\tau} \\ &+\eta^{ii}N_{i0}y_{j}\left(\gamma_{k}y_{l}-y_{k}\gamma_{l}\right)\frac{\gamma^{2}(|+|-\tau)}{\sigma^{2}\tau}-\eta^{ii}N_{i0}y_{j}\left(N_{k0}y_{l}-N_{k0}y_{l}\right)\frac{2\gamma^{6}}{\sigma^{2}\tau} \\ &+\eta^{ii}N_{i0}y_{j}\left(y_{k}N_{0l}-N_{0k}y_{l}\right)\frac{\gamma^{2}(|+|-\tau)}{\sigma^{2}\tau}-\eta^{ii}N_{i0}y_{j}\left(N_{k}y_{l}-N_{k0}y_{l}\right)\frac{2\gamma^{6}}{\sigma^{2}\tau} \\ &+\eta^{ii}N_{i0}(y_{j}(y_{k}-N_{j}y_{k}y_{l})\frac{\gamma^{2}(|+|||^{2}}{\sigma^{2}\tau}+\gamma^{4})+\eta^{ii}N_{i0}N_{0j}\left(y_{k}N_{l}0-N_{k0}y_{l}\right)\frac{2\gamma^{6}}{\sigma^{2}\tau} \\ &+\eta^{ii}N_{ii}y_{j}\left(N_{k}y_{l}-N_{k}y_{l}\right)\frac{\gamma^{2}(|+||^{2}}{\sigma^{2}\tau}+\gamma^{4})+\eta^{ii}N_{i0}N_{0j}\left(y_{k}N_{l}0-N_{k0}y_{l}\right)\frac{2\gamma^{6}}{\sigma^{2}\tau} \\ &+\eta^{ii}N_{ii}y_{j}\left(N_{k}y_{l}-y_{k}y_{l}\right)\left[\left(\gamma^{T}N_{r}\gamma^{2}_{d}-\gamma^{T}\frac{2||\mathbf{y}||^{2}}{\sigma^{2}}-\gamma^{2}\right) \\ &+\eta^{ii}N_{ii}\left(y_{j}y_{l}-\eta_{i}y_{k}\right\right)\left[\left(\gamma^{T}N_{r}\gamma^{2}_{d}-\gamma^{2}-\gamma^{T}\frac{2||\mathbf{y}||^{2}}{\sigma^{2}\tau}\right) \\ &-\eta^{ii}\left(N_{ik}y_{j}y_{l}-N_{il}y_{j}y_{k}\right)\left[\varepsilon^{2}_{d}-\frac{2\gamma^{2}}{\sigma^{2}}-\frac{2\gamma^{3}}{\sigma^{2}}-\left(\varphi\left||\mathbf{y}||^{2}+\omega+\varepsilon\right\right)\frac{2\gamma^{4}}{\sigma^{4}\tau}\right] \\ &+\eta^{ii}\left(N_{ik}y_{j}y_{l}-N_{il}y_{j}\gamma_{k}\right)\left[\varepsilon^{2}_{d}-\frac{2\gamma^{3}}{\sigma^{2}}-\frac{2\gamma^{3}||\mathbf{y}||^{2}}{\sigma^{2}\tau}\right) \\ &-\eta^{ii}\left(N_{ik}y_{j}y_{l}-N_{il}y_{j}y_{k}\right)\left(\omega^{2}_{d}-\frac{2\gamma^{3}}{\sigma^{2}}-\frac{2\gamma^{3}||\mathbf{y}||^{2}}{\sigma^{2}\tau}\right) \\ &-\eta^{ii}\left(N_{ik}y_{j}y_{l}-N_{il}y_{j}y_{k}\right)\left]\varepsilon^{2}_{d}-\frac{2\gamma^{3}}{\sigma^{2}}-\frac{2\gamma^{3}||\mathbf{y}||^{2}}{\sigma^{2}\tau}\right) \\ &+\eta^{ii}\left(N_{ik}y_{j}y_{l}-N_{il}y_{j}y_{k}\right)\left(\frac{2\gamma^{2}}{\sigma^{2}}-\frac{\gamma^{3}||\mathbf{y}||^{2}}{\sigma^{2}\tau}\right) \\ &+\eta^{ii}\left(N_{ik}y_{j}y_{l}-N_{il}y_{j}y_{k}\right)\left(\frac{2\gamma^{2}}{\sigma^{2}}-\frac{2\gamma^{3}}{\sigma^{3}}\left||\mathbf{y}||^{2}\tau}\right) \\ &-\eta^{ii}\left(N_{ik}y_{j}y_{l}-N_{il}y_{j}y_{k}\right)\left(\frac{2\gamma^{2}}{\sigma^{2}}-\frac{\gamma^{3}||\mathbf{y}||^{2}}{\sigma^{2}\tau}-\tau}\right) \\ &-\eta^{ii}\left(N_{ik}y_{j}y_{l}-$$

$$\begin{split} &+\varepsilon \frac{4\gamma}{\tau} + \Re_4 \frac{\gamma^2}{\sigma} - \frac{4\gamma^4 ||y||^2}{\sigma^2 \tau} \Big] \\ &+ \left(\delta_k^i y_j y_l - \delta_l^i y_j y_k \right) \left(\gamma^\tau \gamma_r \right) \left[\left(\varphi ||y||^2 + \varepsilon \right) \omega \frac{\gamma^2}{\sigma} + \frac{2\gamma^2 ||y||^2}{\sigma^2} \right] \\ &+ \left(\delta_k^i y_j \gamma_l - \delta_l^i y_j \gamma_k \right) \gamma_0 \left\{ \frac{2\gamma}{\sigma} \left[\Re_2 - \left(\mathcal{A} ||y||^2 + 2\mathcal{B} + \mathcal{C} \right) \right] \\ &- \frac{\gamma^2}{\sigma^2} \varepsilon \left(4\varphi ||y||^2 + \omega + 4\varepsilon + \alpha + 1 \right) - \frac{\gamma^2 ||y|^2}{\sigma^2 \tau^2} \right\} + \left(\delta_k^i y_j \gamma_l - \delta_l^i y_j \gamma_k \right) N_{00} \frac{\gamma^3}{\sigma^2 \tau} \\ &+ \left(\delta_k^i y_j N_{00} - \delta_l^i y_j N_{k0} \right) \left(\varepsilon \frac{\varepsilon \gamma^4}{\sigma \tau} + N_{00} \frac{\delta \gamma^4}{\sigma^2 \tau} - \gamma_0 \frac{4\gamma^5 ||y||^2}{\sigma^2 \tau} \right) \\ &- \left(\delta_k^i y_j N_{01} - \delta_l^i y_j N_{k0} \right) \left(\gamma (\gamma \cdot \frac{2\beta ||y|^2}{\sigma^2 \tau} - \alpha \frac{2\gamma^2 ||y||^2}{\sigma^2 \tau} \right) + \frac{2\gamma^2 ||y|^2}{\sigma^2 \tau} N_{00} \right] \\ &+ \left(\delta_k^i y_j N_{l} - \delta_l^i y_j N_k \right) \left(\gamma (\gamma \cdot \frac{2\beta ||y|}{\sigma \tau} - \gamma \cdot \sigma \frac{2\beta ||y|^2}{\sigma^2 \tau} \right) + \frac{2\gamma^2 ||y|^2}{\sigma^2 \tau} \right) \\ &- \left(\delta_k^i y_j N_{l} - \delta_l^i y_j N_k \right) \left(\gamma (\gamma \cdot \frac{2\beta ||y|}{\sigma^2 \tau} - \gamma \cdot \gamma \frac{2\beta ||y|^2}{\sigma^2 \tau} \right) \\ &- \left(\delta_k^i N_j \eta_{l} - \delta_l^i N_j \eta_{l} N_k \right) \left(\gamma (\gamma \cdot \frac{2\beta ||y|}{\sigma^2 \tau} - \gamma \cdot \gamma \frac{2\beta ||y|^2}{\sigma^2 \tau} \right) + N_{00} \frac{2\gamma^6}{\sigma^2 \tau} \right) \\ &- \left(\delta_k^i N_j \eta_{l} - \delta_l^i N_j \eta_{l} N_k \right) \left(\gamma (\gamma \cdot \frac{2\beta ||y|}{\sigma^2 \tau} - \varepsilon \frac{2\gamma^4 ||y|^2}{\sigma^2 \tau} \right) \\ &- \left(\delta_k^i N_j \eta_{l} - \delta_l^i N_j \eta_{l} N_k \right) \left(\gamma (\gamma \cdot \frac{2\beta ||y|^2}{\sigma^2 \tau} - \varepsilon \frac{2\gamma^4 ||y|^2}{\sigma^2 \tau} \right) \\ &- \left(\delta_k^i N_j \eta_{l} - \delta_l^i N_j \eta_{l} N_k \right) \left(\gamma (\gamma \cdot \frac{2\gamma ||y|^2}{\sigma^2 \tau} - \varepsilon \frac{2\gamma^4}{\sigma^2 \tau} \right) \\ &+ \left(\delta_k^i N_j \eta_{l} - \delta_l^i N_{j} N_k \right) \left(\gamma (\gamma \cdot (\omega \frac{2\gamma}{\sigma^2 \tau} - \sigma \frac{2\gamma^4}{\sigma^2 \tau} + N_{00} \frac{2\gamma^2 ||y||^2}{\sigma^2 \tau} \right) \\ \\ &- \left(\delta_k^i N_j \eta_{l} - \delta_l^i N_{j} N_k \right) \left(\gamma (\omega \cdot \frac{2\gamma^2}{\sigma^2 \tau} - \varepsilon \frac{2\gamma^4}{\sigma^2 \tau} \right) \\ &- \left(\delta_k^i N_j \eta_{l} - \delta_l^i N_{j} \eta_{l} N_k \right) \left(\gamma (\omega \cdot \frac{2\gamma^2}{\sigma^2 \tau} - \varepsilon \frac{2\gamma^4}{\sigma^2 \tau} \right) \\ \\ &- \left(\delta_k^i N_j \eta_{l} - \delta_l^i N_{j} N_k \right) \left(\gamma (\omega \cdot \frac{2\gamma^2}{\sigma^2 \tau} - \varepsilon \frac{2\gamma^4 ||y|^2}{\sigma^2 \tau} \right) \\ \\ &- \left(\delta_k^i N_j \eta_{l} - \delta_l^i N_{j} N_k \right) \left(\gamma (\omega \cdot \frac{2\gamma^2}{\sigma^2 \tau} - \varepsilon \frac{2\gamma^4 ||y|^2}{\sigma^2 \tau} \right) \\ \\ &- \left(\delta_k^i N_j \eta_{l} - \delta_l^i N_{j} N_k \right) \left(\gamma (\omega \cdot \frac{2\gamma^2}{\sigma^2 \tau} - \varepsilon \frac{2\gamma^4 ||y|^2}{\sigma^2 \tau} \right) \\ \\ &- \left(\delta_k^i N_j \eta_{l} - \delta_l^i N_{j} N_k \right) \left(\frac{\gamma ||y|^2 (2$$

$$\begin{split} &+ \left(\varepsilon + \beta + \varphi \|y\|^2\right) - \left(\omega - \frac{2|y|^2}{2|x|^2}\right) \frac{2\gamma^4}{2\tau^7} + \left(\varphi - \frac{\gamma}{2\tau}\right) \left[1 + \frac{\gamma^2(1-\sigma)}{2x^{3+2}}\right] \\ &- \frac{2\gamma^3}{\sigma^2} - \frac{4\gamma^5(2-3\sigma)}{\sigma^2\tau^2} + \frac{8\varepsilon\gamma^4}{\sigma\tau} - \frac{\gamma^2}{\sigma^2}\right) + \left(\varphi + \frac{2\gamma\gamma}{2\sigma^2}\right) \left(-\frac{\gamma^2}{\sigma}\right) \\ &- \frac{2\gamma^3(2\gamma^2+4+2\tau)}{\sigma^2\tau^2} + \gamma^5||y||^2(9-8\sigma) - \omega \frac{4\gamma^4}{\sigma\tau} \\ &+ \left(\varepsilon + \beta + \varphi \|y\|^2\right) - \left(\omega - \frac{2|y|^2}{2\sigma^2\tau}\right) \frac{2\gamma^4}{\sigma\tau} + \frac{\gamma^3(4\sigma^2-1)}{2\sigma^2\tau^2}\right] \\ &+ y^i y_i (N_0 ky_l - y_k N_0 l) \frac{2\gamma^6(1+\tau-3-4\gamma^2)}{\sigma^2\tau^2} - y^i y_j (N_k 0y_l - y_k N_0 l) \frac{\gamma^6(+21\tau+21+8\gamma^2)}{\sigma^2\tau^2} \\ &+ y^i (N_k jy_l - N_{li} y_k) \frac{2\gamma^4(2\sigma^2+3\gamma^2)}{\sigma\tau^2} + y^i N_{0j} \left[y_k y_l \frac{2\gamma^6(1\tau\sigma-3)}{\sigma^2\tau^2}\right] \\ &+ y^i (N_k jy_l - N_{li} y_k) \frac{2\gamma^4(2\sigma^2+3\gamma^2)}{\sigma\tau^2} + y^i N_{0j} \left[y_k y_l \frac{2\gamma^6(1\tau\sigma-3)}{\sigma^2\tau^2}\right] \\ &+ y^i (N_k jy_l - N_{li} y_k) \frac{2\gamma^4(\tau-3\gamma^2)}{\sigma\tau^2} + y^i N_{0j} \left[y_k y_l \frac{2\gamma^6(1\tau\sigma-3)}{\sigma^2\tau^2}\right] \\ &+ y^i y_j \eta_{kl} \gamma_0 \left\{ \left[8(1-\sigma) + \frac{\sigma-\tau}{\sigma\tau}\right] \frac{\gamma^3(16\sigma^2-1)}{\sigma^2\tau^2} - \frac{4\gamma^7||y||^4}{\sigma^2\tau} + 2\varphi \frac{\gamma^2(2\gamma^2-\tau)}{\sigma\tau} \\ &+ \left(2\omega - \varepsilon + \varphi ||y||^2\right) \frac{2\gamma^4}{\sigma\tau} - \frac{4\gamma^5}{\sigma\tau^2} - \frac{2\gamma^5(1y||^2}{\sigma\tau^2}\right) \\ &+ y^i y_j \eta_{kl} \gamma_0 \left\{ \left[8(1-\sigma) + \frac{\sigma-\tau}{\sigma\tau}\right] \frac{\gamma^3(16\sigma-1)}{\sigma^2\tau^2} + y^i y_j \eta_{kl} N_{00} \frac{\sigma^6(35\sigma-11)}{\sigma^2\tau^2} \\ &+ y^i y_j \delta_{kl} \gamma_0 \frac{4\gamma^3}{\sigma^2} + y^i (\eta_{jk} N_{li} - \eta_{jl} N_{kl}) \\ &+ 2\frac{\gamma^4}{\sigma\tau} - 8\varphi \gamma^2 - \alpha \frac{\gamma^2}{\sigma^2} + \frac{\gamma^3(8\sigma-1)}{\sigma^2\tau} + \frac{2\gamma^2(2\alpha-\gamma)}{\sigma\tau^2} - \frac{2\gamma^5(5\sigma-1)}{\sigma^2\tau^2} \\ &+ y^i (\eta_{jk} N_i - \eta_{jl} N_k) \frac{\gamma^4(2\gamma^2-\sigma)}{\sigma\tau^2} + y^i (\eta_{jk} N_{li} - \eta_{jl} N_{kl}) \frac{\gamma^4(2\gamma^2-\sigma)}{\sigma\tau^2} \\ &+ y^i (\eta_{jk} N_l - \eta_{jl} N_k) \frac{\gamma^4(2\gamma^2-\sigma)}{\sigma\tau^2} + y^3 ||y||^2 (14\tau+25) + \frac{8\varphi \gamma^3 \varepsilon^2 + 2\gamma}{\sigma\tau^2} \\ &+ y^i (\eta_{jk} N_l - \eta_{jl} N_k) \frac{\gamma^4(2\gamma^2-\tau)}{\sigma\tau} + \frac{\gamma^3||y||^2}{\sigma^2\tau^2} - \frac{\gamma^2(\omega+\gamma)||^2}{\sigma\tau^2} \\ &+ y^i (\eta_{jk} N_l - \eta_{jl} N_k) \frac{\gamma^4(2\gamma^2-\tau)}{\sigma\tau^2} + \frac{\gamma^3||y||^2}{\sigma\tau^2} - \frac{\gamma^2(\omega+\gamma)||^2}{\sigma\tau^2} \\ \\ &+ y^i (\eta_{jk} N_l - \eta_{jl} N_k) \frac{\gamma^4(2\gamma^2-\tau)}{\sigma\tau^2} + \frac{\gamma^3||y||^2}{\sigma\tau^2} - \frac{\gamma^2(\omega+\gamma)||^2}{\sigma\tau^2} \\ \\ &+ y^i (\eta_{jk} N_l - \eta_{jl} N_k) \frac{\gamma^4(2\gamma^2-\tau)}{\sigma\tau^2} + \frac{\gamma^3||y||^2}{\sigma\tau^2} - \frac{\gamma^2(\omega+\gamma)||^2}{\sigma\tau^2} \\ \\ &+ y^i (\eta_{jk} N_l - \eta_{jl} N_k) \frac{\gamma^4(2\gamma^2-\tau)}{\sigma\tau^2} + \frac{\gamma^3||y||^2}{\sigma\tau^$$

$$\begin{split} &+ \left(\delta_{k}^{1} N_{j} 0 y_{l} - \delta_{l}^{1} N_{j} y_{k}\right) \left(\frac{2\pi^{4}}{\sigma \tau^{3}} + \left(\delta_{k}^{1} N_{j} y_{l} - \delta_{l}^{1} y_{j} N_{0l}\right) \frac{2\pi^{4}}{\sigma \tau^{2}} \\ &+ \left(\delta_{k}^{1} y_{j} N_{0l} - \delta_{l}^{1} y_{j} N_{0l}\right) \left[\frac{2\pi^{2}}{\sigma \tau^{2}}\right] + \left(\frac{2\pi^{2}}{\sigma \tau^{2}}\right) + \left(\frac{2\pi^{2}}{$$

$$\begin{split} &+y^{i}y^{j}\left[(N_{rk}+N_{kr})N_{l}^{r}-(N_{rl}+N_{lr})N_{k}^{r}\right]\frac{\gamma^{4}(2\gamma^{2}-\tau)}{\sigma^{2}\tau} \\ &+y^{i}y^{j}\left(N_{rk}N_{lr}-N_{kr}N_{rl}\right)\eta^{r\tau}\frac{\gamma^{4}(2\gamma^{2}-\tau)}{\sigma^{2}\tau} \\ &+y^{i}N_{rj}\left(y_{k}N_{rl}-N_{rk}y_{l}\right)\eta^{r\tau}\frac{\gamma^{4}(2\gamma^{2}-\tau)}{\sigma^{2}\tau} \\ &+y^{i}N_{rj}\left(y_{k}N_{rl}-N_{rk}y_{l}\right)\eta^{r\tau}\frac{\gamma^{4}(2\gamma^{2}-\tau)}{\sigma^{2}\tau} \\ &+\eta^{ii}N_{i0}y_{l}\left(N_{ko}-\eta_{ik}N_{l0}\right)\frac{\gamma^{4}}{\sigma^{2}} + \eta^{ii}N_{i0}y_{j}\left(\gamma_{k}y_{l}-y_{k}y_{l}\right)\frac{\gamma^{2}(2\tau+1+\gamma)}{\sigma^{2}\tau} \\ &+\eta^{ii}N_{i0}y_{l}\left(N_{ko}-N_{kl}y\right)\frac{2\gamma^{4}}{\sigma^{2}} + \eta^{ii}N_{i0}y_{j}\left(y_{k}y_{l}-N_{0k}y_{l}\right)\frac{\gamma^{4}(2\gamma^{2}-\tau)}{\sigma^{2}\tau} \\ &+\eta^{ii}N_{i0}\left[\left(N_{jk}+N_{kj}\right)N_{0l}-\left(N_{jl}+N_{lj}\right)N_{0k}\right]\frac{\gamma^{4}(2\gamma^{2}-\tau)}{\sigma^{2}\tau} \\ &+\eta^{ii}N_{i0}\left[\left(N_{jk}+N_{kj}\right)N_{0l}-\left(N_{jl}+N_{lj}\right)N_{0k}\right]\frac{\gamma^{4}(2\gamma^{2}-\tau)}{\sigma^{2}\tau} \\ &+\eta^{ii}N_{i0}\left(\eta_{k}\gamma_{l}-\eta_{j}\eta_{k}y_{l}\right)\frac{\gamma^{3}|y|^{2}}{\sigma^{2}} \\ &+\eta^{ii}(N_{i0}\left(\eta_{k}\gamma_{l}-\eta_{j}\eta_{k}y_{l}\right)\frac{\gamma^{3}|y|^{2}}{\sigma^{2}} \\ &+\eta^{ii}(N_{i0}\left(\eta_{k}\gamma_{l}-\eta_{j}\eta_{k}y_{l}\right)\frac{\gamma^{3}|y|^{2}}{\sigma^{2}} \\ &+\eta^{ii}(N_{i0}\left(\eta_{k}y_{l}-\eta_{il}y_{k}\right)\frac{\gamma^{3}|y|^{2}}{\sigma^{2}} \\ &+\eta^{ii}(N_{i0}\left(\eta_{k}y_{l}-\eta_{il}y_{k}y\right)\frac{\gamma^{3}|y|^{2}}{\sigma^{2}} \\ &+\eta^{ii}(N_{ik}y_{l}-\eta_{il}y_{k}y_{l})\frac{\gamma^{3}|y|^{2}}{\sigma^{2}} \\ &+\eta^{ii}\left(N_{ik}y_{l}-\eta_{il}y_{k}y_{l}\right)\gamma_{j}\frac{\gamma^{3}|y|^{2}}{\sigma^{2}} \\ &-\eta^{ii}\left(N_{ik}y_{l}-\eta_{il}y_{k}y_{l}\right)\gamma_{j}\frac{\gamma^{4}|y|^{2}}{\sigma^{2}} \\ &-\eta^{ii}\left(N_{ik}y_{l}-\eta_{il}y_{k}y_{l}\right)\gamma_{j}\frac{\gamma^{4}|y|^{2}}{\sigma^{2}} \\ &-\eta^{ii}\left(N_{ik}y_{l}-\eta_{il}y_{k}y_{l}\right)\left(\gamma_{l}\gamma^{2}|y|^{2}\right) \\ &+\eta^{ii}\left(\eta_{jk}y_{l}-\eta_{jl}y_{k}y_{l}\right)\left(\gamma_{l}\gamma^{2}\frac{\gamma^{4}|y|^{2}}{\sigma^{2}}} \\ &-\eta^{ir}\gamma^{4}}{\sigma^{2}} \\ &-\eta^{ii}N_{ir}\left(\eta_{jk}y_{l}-\eta_{jl}y_{k}y_{l}\right)\left(N_{00}\frac{\gamma^{4}(\gamma^{2}-\gamma)}{\sigma^{2}}} \\ &-\eta^{ir}\gamma^{4}}(\gamma_{l}y_{k}y_{l}-N_{i}y_{k}y_{l}\right)\frac{\gamma^{4}}{\sigma^{4}} \\ &+y^{i}\left(\eta_{jk}y_{l}-\eta_{jl}y_{k}y_{l}\right)\left(N_{00}\frac{\gamma^{4}(\gamma^{2}-\gamma)}{\sigma^{2}}} \\ &-\eta^{ii}\left(\eta_{jk}y_{l}-\eta_{jl}y_{k}y_{l}\right)\left(N_{00}\frac{\gamma^{4}(\gamma^{2}-\gamma)}{\sigma^{2}}} \\ &-\eta^{i}\left(\eta_{jk}y_{l}-\eta_{jl}y_{k}y_{l}\right)\left(N_{00}\frac{\gamma^{4}(\gamma^{2}-\gamma)}{\sigma^{4}}} \\ &+y^{i}\left(\eta_{jk}y_{l}-\eta_{jl}y_{k}y_{l}\right)\left(N_{00}\frac{\gamma^{4}(\gamma^{2}-\gamma)}{\sigma^{2}}} \\ &+y^{i}\left(\eta_{jk}y_{l}-\eta_{jl}y_{k}y_{l}\right)\left(N_{00}$$

$$\begin{split} &+\gamma^{i}y_{j}\left(y_{k}\gamma_{l}-\gamma_{k}y_{l}\right)\frac{\left(2\sigma-1\right)\left(6\sigma+1\right)}{2\sigma^{2}\tau}+\gamma^{i}|y_{l}\left(N_{kl}-N_{lk}\right)\frac{2\gamma^{3}||y||^{2}}{\sigma^{2}\tau^{2}}\right)}{\sigma^{2}\tau^{2}} \\ &-\gamma^{i}\left(N_{k}y_{l}y_{l}-N_{l}y_{k}\right)\left(\frac{\gamma^{4}\left(\tau-2\gamma^{2}\right)}{\sigma^{2}\tau}+\gamma^{3}||y||^{2}\left(3\sigma-4\gamma^{2}\right)}{\sigma^{2}\tau^{2}}\right) \\ &+\left(N_{k}^{i}y_{l}y_{l}-N_{k}^{i}y_{l}y_{l}\right)\left(\frac{\gamma^{4}\left(1-\sigma^{2}}{\sigma^{2}\tau^{2}}\right)-\left(N_{l}^{i}y_{l}N_{0k}-N_{k}^{i}y_{l}y_{l}y_{l}\right)\right)\frac{\gamma^{4}}{\sigma^{4}} \\ &+\left(N_{l}^{i}y_{l}y_{k}-N_{k}^{i}y_{l}y_{l}y_{l}\right)\frac{2\gamma^{3}}{\sigma^{2}\tau^{2}}-\left(N_{l}^{i}N_{0}y_{k}-N_{k}^{i}y_{l}y_{l}\right)\right)\frac{\gamma^{4}}{\sigma^{4}\tau^{2}} \\ &+\left(N_{l}^{i}y_{l}y_{k}-N_{k}^{i}y_{l}y_{l}\right)\frac{2\gamma^{3}}{\sigma^{2}\tau^{2}}+\left(N_{l}^{i}y_{l}\gamma_{k}-N_{k}^{i}y_{l}y_{l}\right)\left(\frac{9\gamma^{4}}{\sigma^{4}\tau^{2}}\right) \\ &-\left(N_{l}^{i}\gamma_{l}y_{k}-N_{k}^{i}y_{l}y_{l}\right)\frac{2\gamma^{3}}{\sigma^{2}\tau^{2}}+\left(N_{l}^{i}y_{l}y_{k}-N_{k}^{i}y_{l}y_{l}\right)\left(\frac{9\gamma^{4}}{\sigma^{4}\tau^{2}}\right) \\ &-\left(N_{l}^{i}\gamma_{l}y_{k}-N_{k}^{i}y_{l}y_{l}\right)\frac{2\gamma^{3}}{\sigma^{2}\tau^{2}}+N_{l}^{i}\left(y_{k}\gamma_{l}-\gamma_{k}y_{l}\right)\left(\gamma^{2\gamma^{4}}\right) \\ &-\left(N_{l}^{i}y_{l}y_{k}-N_{k}^{i}y_{l}y_{l}\right)\frac{\gamma^{4}}{\sigma^{2}\tau^{2}}+N_{l}^{i}\left(y_{k}\gamma_{l}-\gamma_{k}y_{l}\right)\left(\gamma^{2\gamma^{4}}\right) \\ &-\left(N_{l}^{i}y_{l}y_{k}-N_{k}^{i}y_{l}y_{l}\right)\frac{\gamma^{4}}{\sigma^{2}\tau^{2}}+N_{l}^{i}\left(y_{k}\gamma_{l}-\gamma_{k}y_{l}y_{l}\right)\gamma^{2\gamma^{4}}\right) \\ &-\left(N_{l}^{i}y_{l}y_{k}-N_{k}^{i}y_{l}y_{l}\right)\gamma_{l}\left(\gamma_{l}y_{k}-N_{k}^{i}y_{l}y_{l}\right)\gamma^{2\gamma^{4}}\right) \\ &-\left(N_{l}^{i}y_{l}y_{k}-N_{k}^{i}y_{l}y_{l}\right)\gamma_{l}\left(\gamma_{l}y_{k}y_{l}-N_{k}y_{l}y_{l}\right)\gamma^{2\gamma^{4}}\right) \\ &-\left(N_{l}^{i}y_{l}y_{k}+N_{k}y_{l}\right)y_{l}\left(N_{l}y_{l}+N_{l}y_{l}y_{l}\right)\frac{\gamma^{4}\left(2\gamma^{2}-\tau\right)}{\sigma^{2}\tau^{2}}}\right) \\ &-\left[N_{l}^{i}\left(N_{l}y_{k}+N_{k}y_{l}\right)N_{l}\left(N_{l}y_{l}+N_{k}y_{l}y_{l}\right)\frac{\gamma^{4}\left(2\gamma^{2}-\tau\right)}{\sigma^{2}\tau^{2}}}\right] \\ &-\left[N_{l}^{i}\left(N_{l}y_{k}+N_{k}y_{l}y_{l}\right)\frac{2\gamma^{5}}{\sigma^{2}\tau^{2}}-N_{0}^{i}y_{l}\left(N_{k}y_{k}y_{l}-N_{k}y_{l}y_{l}\right)\gamma_{l}\left(\gamma^{2}\frac{2\gamma^{5}\left(|y||^{2}}{\sigma^{2}\tau^{2}}-N_{0}^{2}\frac{2\gamma^{5}}{\sigma^{2}}\right)}\right) \\ &+\left(N_{l}y_{l}y_{l}\left(N_{k}y_{l}y_{l}-N_{k}y_{l}\right)\frac{2\gamma^{5}\left(|y||^{2}}{\sigma^{2}\tau^{2}}+N_{0}^{2}\frac{2\gamma^{5}}{\sigma^{2}\tau^{2}}}\right) \\ &+\left(N_{l}y_{l}y_{l}\left(N_{k}y_{l}y_{l}-N_{k}y_{l}\right)\left(\gamma^{2}\frac{2\gamma^{5}\left(|y||^{2}}{\sigma^{2}\tau^{2}}+N_{0}\frac{2\gamma^{5}}{\sigma^{2}\tau^{2}}}\right) \\ &+y^{i}y_{l}\left(N_{k}y$$

$$\begin{split} & \mathcal{C} = \frac{\omega}{\gamma} + \frac{2||y||^2}{\sigma^2} + \frac{2\omega}{\sigma\gamma}, \qquad \mathcal{R}_3 = \frac{\omega}{\gamma} + \frac{2||y||^2}{\sigma^2} + \frac{2\omega}{\sigma\tau} + \frac{\varepsilon\gamma||y||^2}{2\sigma^3} + \omega^2, \\ & \mathcal{R}_4 = \varphi(\varphi ||y||^2 + \omega + \varepsilon) \\ & \text{and} \\ & \mathcal{T}_1 = 2\varphi \frac{\gamma^2(2\gamma^2 - \sigma)}{\sigma\tau} + \frac{2\varphi - \gamma^2}{\sigma} + \frac{(2 - 4\gamma^4)(\varepsilon + \varphi ||y||^2)}{\sigma\tau}, \\ & \mathcal{T}_2 = (\varphi - \frac{\gamma}{2\sigma^3}) \frac{\gamma^2(2\gamma^2 - \sigma)}{\sigma\tau} - \frac{\gamma^2}{\sigma} \left(\varphi + \frac{\gamma}{2\sigma^3}\right) + \varepsilon + \varphi ||y||^2 + \beta - \\ & - \frac{2\gamma^4}{\sigma\tau} \left(\omega - \frac{\gamma||y||^2}{2\sigma^3}\right), \\ & \mathcal{T}_3 = -\frac{2\gamma^4}{\sigma\tau} \left[\omega + 2\varepsilon + (3\varphi + \alpha) ||y||^2\right] + 2\varphi \frac{\gamma^2(2\gamma^2 - \sigma)}{\sigma\tau} + \frac{4\gamma^4}{\sigma\tau} \left(\omega + \varphi ||y||^2\right), \\ & \mathcal{T}_4 = \frac{4\gamma^4}{\sigma\tau} - \alpha \frac{\gamma^2}{\sigma} - 8\varphi \gamma^2 + \frac{4\gamma^2}{\sigma\tau^2} - 2\frac{\gamma^4}{\sigma\tau} \left[(\alpha + 2\varphi) ||y||^2 + \omega + \varepsilon\right], \\ & \mathcal{T}_5 = \frac{2\gamma^4}{\sigma^2} - \frac{4\gamma^6}{\sigma^2\tau} - \frac{12\gamma^6}{\sigma\tau^2}, \\ & \mathcal{T}_6 = 4\varepsilon \frac{\gamma^4}{\sigma\tau} - \frac{2\gamma^4 ||y||^2}{\sigma\tau} \left(\varepsilon + \varphi ||y||^2\right) - \frac{\gamma^2}{\sigma} \left(3\varepsilon + \varphi ||y||^2\right) + \frac{2\gamma^3 ||y||^2}{\sigma\tau}, \\ & \mathcal{T}_7 = (\omega + \varepsilon) \frac{\gamma^2(2\gamma^2 - \sigma)}{\sigma\tau}, \\ & \mathcal{T}_8 = -\frac{\gamma^2}{\sigma} \left(\varepsilon + \varphi ||y||^2\right) - \frac{2\gamma^3 ||y||^4}{\sigma^2}, \\ & \mathcal{T}_9 = \frac{4\gamma^3}{\sigma^2} + \frac{\gamma^2}{\sigma} \left(2\varphi + \frac{\gamma}{\sigma^3}\right), \\ & \mathcal{T}_{10} = -\frac{\gamma^2}{\sigma} \varphi ||y||^2 + \varepsilon + \frac{2\gamma}{\sigma^2} \left(\gamma^2 ||y||^4 - \sigma\right). \end{split}$$

Remark 4. In our opinion, the very complicated form of the d-curvatures, produced by the Lagrangian (3), suggests that the Miron-Anastasiei geometrical theory applied for our anisotropic optical Lagrangian does not offer interesting geometrical aspects for the physical properties of the anisotropic medium. Consequently, the physical interpretations of the geometrical results of this paper still remain an open problem. It follows that another geometrical approach for the anisotropic optical Lagrangian (3) is required.

Acknowledgements. We have benefitted from the criticisms of referees of Analele Științifice ale Universității Ovidius Constanța upon previous variants of our paper.

References

- M. Anastasiei, H. Shimada, *The Beil metrics associated to a Finsler space*, Balkan J. Geom. Appl., 3, no. 2, (1998), 1-16.
- [2] V. Balan, Synge-Beil and Riemann-Jacobi jet structures with applications to physics, Int. J. Math. Math. Sci., 2003, No. 27, (2003), 1693-1702.
- [3] V. Balan, M. Neagu, Jet Single-Time Lagrange Geometry and Its Applications, John Wiley & Sons, Inc., Hoboken, New Jersey, 2011.
- [4] R.G. Beil, Comparison of unified field theories, Tensor N.S., 56 (1995), 175-183.

- [5] L.D. Landau, E.M. Lifshitz, Physique Théoretique. 1. Mécanique. 2. Théorie des Champs (in French), Éditions Mir, Moscou, 1982, 1989.
- [6] R. Miron, M. Anastasiei, The Geometry of Lagrange Spaces: Theory and Applications, Kluwer Academic Publishers, Dordrecht, 1994.
- [7] R. Miron, T. Kawaguchi, *Relativistic geometrical optics*, Int. J. Theor. Phys., 30, no. 11 (1991), 1521-1543.
- [8] M. Neagu, Riemann-Lagrange geometry for relativistic multi-time optics, Seminarul de Mecanică., Univ. de Vest Timişoara, Romania 87 (2004), 1-16.
- [9] M. Neagu, A. Oana, An anisotropic geometrical approach for extended relativistic dynamics, Bulletin of the Transilvania University of Braşov, 9(58), no. 1, (2016), Series III: Mathematics, Informatics, Physics, 91-96.
- [10] M. Neagu, A. Oana, V.M. Red'kov, An anisotropic geometrical approach for non-relativistic extended dynamics, Ricerche Mat., 62, no. 2, (2013), 323–340, DOI 10.1007/s11587-013-0154-8.
- [11] A. Szász, Beil metrics in complex Finsler geometry, Balkan J. of Geom. Appl., 20, no. 2, (2015), 72-83.

Mircea Neagu and Alexandru Oană